Image analysis of bioparticles accumulation and diamagnetic alignment in high-gradient magnetic field

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Abstract. Magnetic properties of biological particles are measured in high-gradient magnetic separation (HGMS) analysis, revealing the concentrating process of nucleoprotein particles, ferritin, red blood cells, and eggs. A magnetic force acting on micrometer and submicrometer biological particles having diamagnetic or paramagnetic susceptibility with respect to the solution causes their movement and accumulation in gradient magnetic fields dependent on the values of the magnetic moments. The methods developed enable us to obtain the magnetic moments values of single particles and their assembly directly from magnetic separation and image analyses without assuming the detection of sizes. Our precision methods for the measurement of the capture traveling (magnetic diffusion) time and the accumulation (magnetic sedimentation) radius in HGMS show that it is really possible to determine the weak diamagnetic shifts of magnetic susceptibility up to $0.7 \times 10^{-10}$ (SI units). HGMS analysis of the concentrating process of nucleoprotein granules (microcells, DNA granules, or nucleosome core particles) with polarization microscopy reveals phase transitions for DNA in granules, and separation accumulation of particles enables the determination of the diamagnetic susceptibility and anisotropy properties. Magnetic concentration effects always occur in living systems because micrometer-located gradient magnetic fields inside an organism are strong enough to cause drifts of cellular complexes and organelles of micrometer and submicrometer sizes. We report the appearance of superparamagnetic contamination inside developing shrimp eggs. In the developing shrimp eggs, ferritin aggregates are observed under weak gradient magnetic fields and diaparaferromagnetic changes are detected. A significant interruption of egg development is revealed in such fields.

Keywords: magnetic separation of bioparticles; magnetic susceptibility; magnetic moments; DNA granules; ferritin; red blood cells; nucleosome core particles.

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1 Introduction

Important achievements in magnetobiology and biomedical physics during last decades have changed our understanding of the biophysical mechanisms of processes in living systems. Recently, the particle-separation biotechnology was developed and has become an important branch of immunodiagnostics. Magnetic separation and magnetophoresis had been applied mainly in cellular assays for removing in gradient magnetic fields of labeled particles, immunomagnetic beads, and targeted antigens.1 For weakly magnetic particles, diamagnetic and paramagnetic ones, the high gradient magnetic separation was developed for industrial purification.2–6 and the technological analysis of particles and cells.7–9 The other important application of high magnetic fields is an alignment method for macromolecules and biomicroparticles, combined with other techniques,7 which enables the detection of the diamagnetic structure of anisotropic particles. In both directions, the physical-chemical backgrounds of methods are connected with the concentration processes of biomacromolecules, which are the subject of research also using the methods of colloidal chemistry, and optical methods, namely, polarization microscopy and light scattering.10–16

In this paper, the high-gradient magnetic separation (HGMS) method shows the concentration process of macromolecules of nucleoprotein, ferritin, and blood cells, and the approach developed with polarization microscopy reveals phase transitions for DNA protein granules. Nucleoprotein granules (DNA granules called Bonghan granules) investigated in this study are suggested as messengers in the acupuncture meridian system and sensitive biosensors working on the basis of optical changes under physical acts.17,18 The separation accumulation of biological particles in HGMS enables the determination of diamagnetic susceptibility.19–21 The
coexistence of isotropic, cholesteric, and smectic phases introduced for colloidal suspensions studies, appeared in our study as the coexistence of a sign variety in magnetic susceptibilities. The protein affinity adsorption and the nucleoprotein alignment during a magnetic aggregation in a high-gradient magnetic field are analyzed.

Magnetic drifts, shifts, and vibrations of micro-objects such as cells, organelles, and cellular depositions with respect to other always occur in an organism under even weak magnetic fields, which penetrate inside the body almost without distortion. In magnetotherapy treatments, alternative gradient magnetic fields (with value $H_0 \times \text{grad}H_0$ of the order of $10^4$ to $10^5$ A/m$^2$/m$^3$) enforce ferritin cores, hemosiderin, etc. to aggregate inside the human body followed by vibrations, increasing the oxidative exchange. Ferritin inside mitochondria, hemosiderin bound with macrophages, FeS clusters bound with mitochondria or sideroblasts—all these complexes may possibly vibrate and drift under the magnetotherapy procedure. Measurements were performed for magnetic drifts of ferritin near to hemosiderin depositions will be useful for modeling magnetotherapy.

2 Methods

2.1 Observation System

To observe the concentration phase transitions in cells and nucleoprotein particles (granules) and their aggregates in a high-gradient magnetic field, ferromagnetic wires were positioned in the visual field of the inverted and reflected microscopes supplied with a CCD camera. The external field intensity was $H=8 \times 10^5$ A/m, with a gradient value on the wire surface $H \times \text{grad}H=3.2 \times 10^{12}$ A$^2$/m$^3$. The HGMS observation chamber was made of a very thin rectangular glass tube (0.1 x 3 mm) or plastics cuvette (10 x 5 mm). A very thin capillary was used for the observation of blood cell sedimentation in a dense medium such as diluted whole blood. In the sample hermetic cell, a ferromagnetic nickel wire with a radius of 10, 15, 50, 163, or 315 μm was mounted in the axial configuration (gravity force parallel to wire axis, and field perpendicular to flow and wire) as Fig. 1(a) demonstrates.

The external steady magnetic field created by the structure of permanent pieces containing the “open domain boundary”
Table 1 Methods of affinity ligand detection.

<table>
<thead>
<tr>
<th>Methods Used</th>
<th>Sorbent Volume (µl)</th>
<th>Protein on Beads (µg/µl)</th>
<th>Cuvette Volume (µl)</th>
<th>Tested Amount (µg)</th>
<th>Tested Concentration (µg/ml)</th>
<th>Assay Time</th>
</tr>
</thead>
<tbody>
<tr>
<td>HGMS capture of single bead</td>
<td>0.1–0.3</td>
<td>0.2–9</td>
<td>10–50</td>
<td>0.02–27</td>
<td>0.02–1.0</td>
<td>0.5 min–0.5 h</td>
</tr>
<tr>
<td>Magnetic particle collection</td>
<td>1–5</td>
<td>0.2–6</td>
<td>10–50</td>
<td>0.2–30</td>
<td>0.2–0.6</td>
<td>0.5 min–0.5 h</td>
</tr>
<tr>
<td>Micro-ELISA</td>
<td>1–10</td>
<td>0.01–5</td>
<td>100</td>
<td>0.01–5</td>
<td>0.1–0.5</td>
<td>1–3 h</td>
</tr>
</tbody>
</table>

is presented in Fig. 1(b). The distribution of magnetic field intensity was simulated according to the equation: 
\[ H(x/a) = 4M_x \ln[\cotan(x/a)] \]
where \( \pi a \) is the length of a single permanent magnet piece.

The movement of single particles was observed using a microscopy system combined with a CCD camera shown in Fig. 1(c). The process of time measurements on successive radial coordinate intervals from the initial position of the particle \( r_{0\phi}(5.5a^2/2a) \) to the end point \( r_{a\phi} = 1.25a \) (where \( a \) is the radius of the wire), enabled to the capture motion parameter—
the traveling time \( T \) on intervals. The viewing field in the vicinity of the wire was calibrated using NI Vision Assistant 7.0 software to provide the particles size dimensions, the luminescence intensity, and a coordinate scale along the particles trajectory.

The magnetic susceptibility of single cells was calculated from experimental dependence of the coordinate function \( F \) via the capture time near a ferromagnetic wire, a saturation magnetization of wires consisted \( M_x = 7700 \) to 8200 G.

The nucleoprotein granules (Bonghan granules, or DNA granules) of a rabbit were extracted by the method described in Ref. 17. Human red blood cells (RBCs) diluted by a factor of 100 by phosphate buffer (PBS) (pH 7.4) were separated by a settling procedure to obtain small cells 3 to 4 \( \mu \)m in diameter. The ferritin particles were horse spleen ferritin commercially produced by “Reanal,” and the buffer used was PBS. Shrimps eggs (from Russian lakes) were incubated in 0.5% NaCl solution for 24 h.

2.2 Affinity Procedures

The successive procedures of the present method were described in Ref. 20 to obtain Sepharose 4B beads (control 1 and control 2) and to immobilize the macromolecular ligands in their meshes, which included the incubation of beads with the ligand samples or the injection of a sample through an affinity minicolumn; the submerging of the loaded beads into the glass cell containing the Ni wire or a narrow gap of magnetic poles; and the obtaining of the magnetic parameters: capture traveling time \( T \) (using HGMS), initial buildup radius \( R_{0\phi} \), capture radius \( R_c \), and saturation accumulation radius \( R_{sat} \) (Refs. 8 and 19–21).

Table 1 demonstrates the used ranges of the biotechnological procedure values used for tests of the macromolecular ligand concentration within the meshes of porous agarose beads.

2.3 Preparation of Agarose Immuno-Beads

Sepharose 4B beads with 4% (volumetric percent) of adsorbent net particles (“Pharmacia”) crashed to diameters of 3 to 40 \( \mu \)m were prepared for the ligand immobilization according to the conventional Br-CN-activated procedure.8,14,20

The adjusted optimization of the magnetic susceptibility values was reached by the shaking of clear beads in 1 mM HCl solution for 5 and 30 min to obtain control 1 and control 2, respectively, at 4 \( ^\circ \)C. The immobilized single macromolecular layer was covalently bound to agarose beads after incubation in saline solution PBS (pH = 7.5). Such ligands were bovine serum albumen (BSA), egg albumen (Alb), human immunoglobulin G (Ig G), and “Serva.” Double layers were presented by antigen (Ag) to antibody (Ab) pairs: IgG and goat antihuman immunoglobulin G or as prepared.8,14 To prevent the nonspecific binding processes, which are less common for an affinity sorbent such as agarose, 0.1 mg/ml BSA and 0.02% Tween 20 were added to the samples. Micro-enzyme-linked immunosorbent assay14 (micro-ELISA) enabled us to determine the content of the linked protein on the beads to calibrate the suggested method. Thus, 100 \( \mu \)l of freshly prepared conjugates of diluted peroxidase-labeled secondary antibodies were added to 100 \( \mu \)l tested suspension of beads in saline buffer on a micro-ELISA plate. After shaking for 15 min, the reaction with 10 \( \mu \)l 0.09 mM peroxidase substrate (o-phenylenediamine) and 30 \( \mu \)l 0.5% \( H_2O_2 \) was run and interrupted by 30 \( \mu \)l of 50% \( H_2SO_4 \) after 5 min. The yellow color of the probes was measured by a micro-ELISA plate reader MR 580 at 405 nm. The absorption measurements of protein concentrations in liquid were carried out with the Hitachi spectrophotometer at 280 nm.

3 Results

The buildup process on the wire was theoretically and experimentally studied for new particles: nucleoprotein granules, ferritin aggregates, and blood cells. The saturation accumulation radius of buildup layer on the wire was independent of particles sizes in an axial configuration and a vertical wire. In the axial case, the cylindrical symmetry imposed parallel to a wire produces equations with analytical solutions for the motion of the particle with the plane of the wire magnetization, as Fig. 1(a) shows. The equations of motion in the \( (r, \phi) \) plane are
\[
\begin{align*}
\frac{dr}{dt} &= -(V_m/a)(K/r_a^5 \pm \cos 2\phi/r_a^3), \\
\frac{d\phi}{dt} &= -(V_m/a)(\sin 2\phi/r_a^3),
\end{align*}
\]

where

\[
V_m = 2\chi M_s H_0 b^2/9\eta a,
\]

\[
K = M_s/2H_0,
\]

\[
r_a = r/a.
\]

(1)

Here (+1) is for para/ferromagnetic capture and (−1) for diamagnetic capture; \(a\) and \(b\) are the wire and particle radii, respectively.

Integrating the second motion Eq. (1) along a trajectory from point \((r_{a_1}, \phi_1)\) to \((r_{a_2}, \phi_2)\) one can get the time \(T\) of particle movement:

\[
T = \int_{\phi_1}^{\phi_2} \frac{dr}{c(\sin 2(\phi - \alpha))},
\]

where \(r_a = 1\), \(\phi_2\) is determined from the equation of trajectory, as in the following.

The equation describing the trajectory obtained from the equations of the motion of Eq. (1) is

\[
r_a^2 = c[\sin 2(\phi - \alpha)],
\]

where \(r_a\) and \(\phi_1\) are determined by the initial position of the particle in the plane. For a given trajectory, all particles passing the same distance interval demonstrated the reverse dependence of their magnetic susceptibility; two trajectories are indicated in Fig. 1(a).

The magnetic susceptibility can be obtained from the equality of the ponder-motive force and the Stokes drag force:

\[
6\pi \eta b v = 1/2(\chi_0 - \chi_f) \cdot \nabla (H \cdot H),
\]

where \(\chi_0\) and \(\chi_f\) are the susceptibility of the particles and the fluid; \(\eta\) is the viscosity coefficient; and \(v\) and \(V\) are the velocity and volume of the particle, respectively. For a cylinder magnetized in the field, the solution of the scalar potential equation by Laplace gives

\[
\chi = 6\pi \eta b a^2 F(r_{a_1}, \phi_1; r_{a_2}, \phi_2)/VM_sH_0T,
\]

where

\[
F = c^2[\sin(\phi_1 + \phi_2 - 4\alpha)\sin(\phi_1 - \phi_2) + (1/2)\sin^2 2\alpha \ln(\tan \phi_1/\tan \phi_2)],
\]

where \(\phi_1\) and \(\phi_2\) are azimuth angle coordinates for initial and ending positions, and \(c\) and \(a\) are found according to Eq. (2). Assuming the azimuthal angle \(\phi\) of the particle with the plane of the wire’s magnetization to be zero, the equations of motion can be integrated using Eq. (2), which can be written as

\[
T = -(1/V_m)(y_2 - y_1),
\]

where \(y = [r_a^4/4 - K r_a^2/2 + K^2 \ln(r_a^2 + K)/2]\); \(r_a\) is the normalized coordinate; \(T\) is the traveling time of the particle from position \(r_{a_1}\) to \(r_{a_2}\); and \(M_s\) is the saturation magnetization of the wire in an applied magnetic field \(H_0\) perpendicular to the axis of the wire having radius \(a\). The relative magnetic susceptibility can be calculated from the magnetic velocity \(V_m\) by

\[
\chi = \chi_0 - \chi_f = 9V_m \eta a/2H_0M_s b^2.
\]

The magnetic moments of single red blood cells in oxygenated, met-hemoglobin, and deoxygenated states of containing hemoglobin can be determined with 10% error, as Table 2 shows. The magnetic movement speeds of red blood cells in oxygenated state near wires are shown in Fig. 2.

<table>
<thead>
<tr>
<th>Particles</th>
<th>(\mu \times 10^{15}) erg/G</th>
</tr>
</thead>
<tbody>
<tr>
<td>OxyHb red blood</td>
<td>−0.57 ± 0.05</td>
</tr>
<tr>
<td>DeoxyHb red blood</td>
<td>22.5 ± 0.8</td>
</tr>
<tr>
<td>MetHb red blood</td>
<td>27.4 ± 0.7</td>
</tr>
<tr>
<td>Nucleoprotein granules</td>
<td>−0.53 ± 0.08 × 10⁻³</td>
</tr>
<tr>
<td>Ferritin aggregates</td>
<td>145.4 ± 5.1</td>
</tr>
<tr>
<td>Ferritin cores</td>
<td>2400 ± 30</td>
</tr>
</tbody>
</table>

*1 erg/G = 10⁻³ Am².

One can estimate the molar magnetic susceptibility of oxy-hemoglobin assuming a contamination of red blood cells −28%, a hemoglobin density of 1.335 g/cm³, and a molecular weight value MW = 64,500: \(\chi_M = -0.036\) cm³/mole, for deoxy-Hb: \(\chi_M = 0.0118\) cm³/mole, for MetHb: \(\chi_M = 0.0132\) cm³/mole. Using the Curie law,

\[
\mu_{eff} = \left(\frac{\chi_M 3kT}{N_A}\right)^{1/2} = 2.83 \mu_B(\chi_M T)^{1/2},
\]

so for high-spin (five unpaired electrons) states of Fe³⁺-hemes in methemoglobin: \(\mu_{eff} = 5.6 ± 0.1\mu_B\), in deoxy-hemoglobin: \(\mu_{eff} = 5.3 ± 0.1\mu_B\).

Figure 2 shows capture time measurements for red blood cells moving along one trajectory to diamagnetic region of the wire. The linear dependencies of the coordinate function \(F\) on time were used to determine the magnetic susceptibility of single cells.
Experimentally measured capture time for immunoparticles moving on the same trajectories are presented in Fig. 3 via the concentration of bound protein BSA, immunoglobulin, and antigen-antibody complexes.

Figure 4 illustrates typical curves of diamagnetic capture of immunobeads with concentration of bound protein as a parameter, capture time via particle diameters, followed an exponential fitting dependence.

The nucleoprotein granules submerged in a high-gradient magnetic field appeared to aggregate under accumulation to the wire (Fig. 5). The fact of saturation accumulation of the build-up process on the wire is illustrated by the curve $R_a(t)$ in Fig. 5, obtained for cylindrical droplet growth (photo in Fig. 5).

The initial equation for the growth of a cylindrical droplet (with mass $m$ and the base area $S$) for time $t$ along the radial coordinate direction is

$$
\frac{\Delta m}{\Delta t} = \frac{\Delta R_a S_p}{\Delta t} = c S \left( \frac{dr}{dt} \right),
$$

where the velocity $dr/dt$ of particles is described by the first equation in Eq. (1) with $\theta=0$. Because the exact Stokes drag force is given by $f_D = \text{const} \cdot \pi r^2 \rho_p v_p^2$, where $\beta$ represents spherical particles and is 1, and $v_p$ is the velocity of the particle, the growth of a cylinder of nucleoprotein granules on wire is described by

$$
\frac{\Delta R_a}{\Delta t} = \frac{c}{\rho_m V_m} \left( \frac{1}{R_a} - \frac{1}{R_a^3} + \frac{K}{R_a^5} \right)^{1/\beta}.
$$

With $K \ll 1$, we get

$$
\frac{dR_a}{dr} = \frac{c V_m}{\rho_m R_a^{3+1}}, \quad \text{so}, \quad R_a^{3+1} = \frac{\rho_m t}{c V_m} + 1.
$$

The obtained parameter value $\beta=1.5$ reflects the nonspherical shape of the DNA granules.
The good sigmoid fitting shown in Fig. 5 could mean the cooperative behavior of granules having induced diamagnetic moments and attracting each other, but the polynomial fitting was the best.

The sensing principle for the biosensor was defined by Cullum and Vo-Dinh\textsuperscript{15} as a device that consists of a receptor and a transducer, placed in the sample environment, where the target molecules, particles, or cells bind to the receptors. In the result of interaction between the analyte and the receptor, the effect can be measured by the transducer, which transforms this effect into the detected signal. The magnetic biosensor developed in this study is based on the magnetic attraction of micrometer and submicrometer bioparticles to magnetized ferromagnetic wire and the accumulation of bioparticles in the detected region.

For high-gradient magnetic separation of submicrometer particles of nucleoprotein and ferritin we detected the growth with the saturation kinetics of the flocculating center of accumulated particles.\textsuperscript{21} The process can be described through an equation of continuity for particle flow in a gradient magnetic field, and the concentration distribution follows an exponential dependence, the argument for which contains the coordinate-independent coefficient $G_0$, which is proportional to the average effective magnetic moment $\mu_p$ of separated submicron particles.

The buildup radius for submicrometer particles can be obtained from equation of continuity

$$\frac{dc}{dr} = \frac{1}{r^2} \frac{G}{r} \frac{\partial (Gc)}{\partial r},$$

and for a steady state, $dc/dr=0$, solution is

$$c = c_0 \exp \left[ - \int_{R}^{r_a} \frac{G}{r} g(x) dx \right], \quad c = c_{\text{sat}},$$

where $c_0$ is the initial particle volume concentration, and assuming $G = G_0 g(x)$, yields for submicrometer particles:

$$2G_0 = -4\pi \chi(M_H) b^3/3kT = \mu_p M/JkT.$$ \textsuperscript{12}

3.1 Polarization Microscopy Study of Magnetic Accumulated Structures

Phase changes were proved by measurements in polarization microscopy studies using CCD camera and LabView 7.0 software. The phase transitions of particles that accumulated on the wire were detected for nucleoprotein granules, but we could not determine the brightness at the same level in polarized light for ferritin and red blood cells.
We obtained completely new method to detect isotropic-cholesteric-smectic phase transitions recording the orientation of DNA granule aggregation in a high-gradient magnetic field.

The photos in Figs. 6a–6i demonstrate the gradual growth of liquid crystalline aggregate of nucleoprotein granules in a high-gradient magnetic field in a suspension and the same aggregate in a dried state. The transformation from an ellipsoid form to a cylindrical form and then to a conic form was completely according to theoretical and experimental research of the “Web-DNA-Collection group,” observing isotropic-nematic-cholesteric-smectic transitions in higher density crystalline states of DNA solutions.

Measurements were carried out to determine the distribution of nucleoprotein and ferritin submicrometer particles along the radial coordinate. Images of the buildup process on wire were treated with LabView 7 software to get $G_0$, as Figs. 5–7 show. Using the luminance distribution values obtained along a radial coordinate $r_a$ direction near the wire, one can use exponential decay fitting to get the parameter $G_0$, which is independent of the wire radius and reflects a logarithmic relation between the saturation concentration $I_s$ of accumulated particles and the initial concentration $I_0$ of them for $r_a=1$:

$$I_s(r_a) = I_0 \exp\left[2 \pi \chi b^3 MH_0/3kT(1/r_a^2 - K/r_a^4)\right].$$

Figure 7 shows the luminance intensity (proportional to concentration) distribution along the radial coordinate direction with exponential fitting. One can obtain magnetic moment values for particles using Eq. (4) and get $G_0$ from exponential fitting, as $\ln(I/I_0)$ for $r_a=1$, $K=0.5$. For the parameters used in the experimental setting, we obtained for the magnetic moment of the particles $\mu_p = \chi VH_0$ in a magnetic field intensity of $8 \times 10^4$ A/m:

$$I_s(r_a) = I_0 \exp\left[2 \pi \chi b^3 MH_0/3kT(1/r_a^2 - K/r_a^4)\right].$$

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For submicron granules with diameters of 0.3 to 0.5 \( \mu \text{m} \), one can estimate the magnetic susceptibility for the available \( G_0 \) interval (0.98 to 2.7) as \( \chi_{\text{gran}} = -0.81 \) to \( 2.53 \times 10^{-8} \) (CGS volumetric units), where the lower value is for \( G_0 \sim 0.98 \).

The magnetic moments of bioparticles are shown in Table 2. They have a large spread of values, and change sign and magnitude considerably between diamagnetic and superparamagnetic ferritin cores form hemosiderin conditions.

Our light reflection and polarization microscopy studies have shown that granules in aggregates have persistence lengths of 5 to 30 \( \mu \text{m} \), while bacteriophage aggregate is 2.2 \( \mu \text{m} \).

The aggregation of granules growing in the manner shown in Fig. 6 was irreversible, as confirmed by observation under strong 1600× magnification with an inverted microscope, as Figs. 8(a) and 8(b) show for the dried granule aggregate in polarized and depolarized light, in the crystalline structure.

The strong mechanical perturbations were applied to force all colloidal particles to vibrate vigorously, but the aggregate obtained [Figs. 6(f)–6(i)] remained accumulated and tended to orient along the external magnetic field being far from a wire. Typically, a permanent long nucleus axis was seen (Fig. 8) as evidence of DNA-continuum formation.

Figure 9 shows the image analysis of a single first right granule of a chain depicted in Fig. 8(a). The image analysis software enabled us to detect a single DNA granule profile and dark diffraction bands, with a granule size of 1.5 \( \mu \text{m} \), as Fig. 9 demonstrates. Diffraction was detected only transverse to the optical axis of the crystalline structure from granules [see Fig. 9(a)]. Along the axis, the bands were absent [see Fig. 9(b)]. Image analysis enabled the recording of a microparticle profile near the wire shown in Fig. 10.

Threadlike structures (collagen tubes) containing DNA granules conferred diamagnetic anisotropic properties to the crystalline optical anisotropic aggregates inside them. Such structures were observed in HGMS attached to the magnetized wire in a diamagnetic area perpendicularly to the wire and the external field direction (Fig. 11).

Based on the observed processes and reported models,\(^{17,18}\) we concluded that the mechanism of light penetration into the threadlike structures in acupuncture points could be such as is shown in the schematic of Fig. 12.

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**Fig. 11** CCD photo of a threadlike tube (Bonghan duct) containing DNA granules in a gradient magnetic field obtained by a polarization microscope. The duct contains the crystalline aggregates of aligned granules.

\[
\mu_p = 0.53 \times 10^{-18} G_0 \times 10^{-3} \text{ (SI units).} \quad (14)
\]

For submicron granules with diameters of 0.3 to 0.5 \( \mu \text{m} \), one can estimate the magnetic susceptibility for the available \( G_0 \) interval (0.98 to 2.7) as \( \chi_{\text{gran}} = -0.81 \) to \( 2.53 \times 10^{-8} \) (CGS volumetric units), where the lower value is for \( G_0 \sim 0.98 \).

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**Fig. 12** Influence of gradient magnetic field (pressure, vibrations, stress, an apoptosis) on granule crystalline phase forming in a threadlike tube [Bonghan (BH) duct]. Areas of crystals arise in ducts due to the magnetic alignment in a high-gradient magnetic field and remain there irreversibly, as previous long-time experiments have shown.

**Fig. 13** (a) Phylum Arthropoda, brine shrimp eggs: *Artemia salina*. Developing eggs of 100-500-\( \mu \text{m} \) size were attracted to a magnetized 625-\( \mu \text{m} \) ferromagnetic wire, \( H_0 \) is 0.05 T, and (b) although repulsed from a magnetized wire when eggs are in earlier stages or when empty, the new shrimps were caught by the gradient magnetic field.
3.2 Ferritin Effects in Developing Artemia Shrimp Eggs

We obtained new data on the magnetic properties of developing eggs in the initial stage of embryos, the developing eggs were found to have a diamagnetic effect with respect to the solution that was revealed in a gradient magnetic field near a magnetized ferromagnetic wire surface, while in the late stages of shrimp eggs, the development changed to paramagnetic eggs [Fig. 13(a)]. Because iron-deficient eggs are restricted in their development, we proposed that magnetic transformation of developing cysts might be due to the change at the amounts of hemoglobin and ferritin amounts during development of the embryo into a nauplius. When the cyst crashed, the nauplius head was first seen, then a new development of the embryo into a nauplius. When the cyst was turned off, while apoferritin protein of ferritin crashed, the nauplius head was first seen, then a new development of the embryo into a nauplius. When the cyst was turned off, while apoferritin protein of ferritin crashed, the nauplius head was first seen, then a new development of the embryo into a nauplius.

3.3 Concentration Effects of Ferritin and Hemosiderin

Figure 14 demonstrates ferritin molecules in gradient field of order $10^5$ to $10^6$ A²/m³. Ferritin cores (hemosiderin) formed a dense layer of about 50 μm. The magnetized in external field and attracted ferritin molecules and aggregates in our model, imitating the depositions of organisms. The ferritin cores were separated by a black layer on the wire surface [Fig. 14(b)] and remained on the wire after the external field was turned off, while apoferritin (protein of ferritin) aggregates diffused immediately. The theoretical curve of ferritin deposition on the magnetized wire would differ from the experimental one, mainly due to the high magnetization of ferritin cores (hemosiderin) collected close to the wire surface.

3.4 Simulation of Deposition in Gradient Fields

A Langevin equation for magnetization of noninteracting superparamagnetic particles is

$$M = VnM_{sp}\left[\coth(VM_{sp}H/kT - kT/VM_{sp}H_0)\right],$$

where $M_{sp}$ is the saturation magnetization of the particle, and $V$ and $n$ are the particle volume and concentrations. Magnetization of superparamagnetic particles such as ferritin cores or FeS clusters of mitochondria (sideroblasts) given by integrating of the Langevin equation follows the function

$$M = VnM_{sp}\int_{\xi} L(\xi)f(m)dm,$$

where $g$ is the number of particles in the aggregate or cluster; $m = gM_{sp}$ is the magnetic moment of the aggregate; $\xi = VH_0M_{sp}/kT$; and $f(m)$ is a function of the distribution that depends on the magnetic polarization $\Delta \alpha$ of a particle due to the aggregation:

$$f(m) = f_0 \exp\left(\frac{\Delta \alpha}{2kTgH_0^2}\right).$$

In the case of particles interacting under their aggregation or clustering, the magnetization value became proportional to $g^2H_0^2$ in weak magnetic fields. A simulation of the deposition of aggregates of super-paramagnetic particles near the magnetized collector [for example, ferritin molecules near a magnetite Fe₃O₄ crystal rod of cylindrical form (radius $a$)] was calculated based on the equation $C_r(x) = C_0 \exp\left(2π χb^3MH_0/3kT \left[1/x^2 + K/x^4\right]\right)$, where $C_r$ is the concentration obtained as a function of $x = r/a$, the normalized distance from the cylinder axis; and $C_0$ is the particle volume concentration for $r = 1$. The pattern obtained for weak fields is demonstrated in Fig. 15, which shows dependence on concentration for ferritin molecules attracted to a...
cylinder in weak fields. A 3% increase in $C$ took place on the surface of the cylinder. For particles of 10-nm sizes with ferritin cores, $M_s = 5100$ G, the equality of the $kT$ value ($0.4 \times 10^{-20}$ J) and magnetic energy $MH_0V$ takes place for $H_0 = 0.25$ Oe = 20 A/m. Since the earth magnetic field is 0.5 Oe ($40 \times 10^{-11}$ CGS units), the steady component, note that the magnetization in the permanent field of the earth for Fe$_3$O$_4$ particles can be significant in living organisms. In strong industrial fields of more than 100 to 500 G, living organisms may become “magnetized” when ferritin and hemosiderin crystals in an organism become aligned, creating many ordered local centers inside the body.

4 Discussion

The magnetic concentration of diamagnetic granules and super-paramagnetic ferritin aggregates$^{21,72}$ occur in points of gradient magnetic applications. Then optical transformations

<table>
<thead>
<tr>
<th>Particles and Influences</th>
<th>Magnetic Susceptibility $(10^8$ CGS units)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Nucleoprotein granules (microcells)</td>
<td>$-1.65$ ($-0.81$ to $-2.53$)</td>
</tr>
<tr>
<td>Ferritin</td>
<td>$+1250$ (3.1 to 2400)</td>
</tr>
<tr>
<td>Red blood cells:</td>
<td></td>
</tr>
<tr>
<td>oxygenated in dilute Tyrode saline (pH 7.4)</td>
<td>$-0.84$ ($-0.62$ to $-1.45$)</td>
</tr>
<tr>
<td>glucose 0.002 mg/ml in Tyrode</td>
<td>$-0.83$ ($-0.41$ to $-1.40$)</td>
</tr>
<tr>
<td>albumen BSA 0.8 mM</td>
<td>$-0.62$ (echinocytes over 30 min)</td>
</tr>
<tr>
<td>5% ethanol over 6 min incubation</td>
<td>$-0.73$ ($-0.41$ to $-1.49$)</td>
</tr>
<tr>
<td>1 h 25 min</td>
<td>$-0.48$ ($-0.42$ to $-0.63$)</td>
</tr>
<tr>
<td>3 h 50 min</td>
<td>$0.33$ ($-0.22$ to $-0.44$)</td>
</tr>
<tr>
<td>$H_2O_2$ 15 mM/h over 8-min incubation</td>
<td>$-0.66$ ($-0.52$ to $-0.80$)</td>
</tr>
<tr>
<td>1 h 20 min</td>
<td>$-0.48$ ($-0.35$ to $-0.74$)</td>
</tr>
<tr>
<td>2 h 40 min</td>
<td>$-0.60$ ($-0.50$ to $-0.70$)</td>
</tr>
<tr>
<td>Methylene blue 5 mM</td>
<td>as normal cells (0.5% of methemoglobin)</td>
</tr>
<tr>
<td>5 mM 20-min incubation</td>
<td>$-0.73$ to $+0.97$ 7% methemoglobin</td>
</tr>
<tr>
<td>1 h 40 min incubation</td>
<td>$-0.15$ ($-0.54$ to $+7.84$) 20% methemoglobin</td>
</tr>
<tr>
<td>Ghosts of red blood cells in 0.1% NaNO$_2$</td>
<td>$+2.08$ to $+17.10$ test of hemoglobin</td>
</tr>
<tr>
<td>Deoxygenated red blood cells (95% of deoxyhemoglobin)</td>
<td></td>
</tr>
<tr>
<td>0.1% sodium dithionite</td>
<td>$+10.0$ to $+30.0$</td>
</tr>
<tr>
<td>Methylene containing erythrocytes by 0.01% NaNO$_2$</td>
<td>$+22.0$ to $+35.0$</td>
</tr>
<tr>
<td>Lymphocytes</td>
<td>$-0.37$ to $-0.46$</td>
</tr>
<tr>
<td>Weak UV irradiation</td>
<td>$-0.51$ to $-0.65$</td>
</tr>
<tr>
<td>Strong UV intensity</td>
<td>$+22.0$ to $+34.0$</td>
</tr>
<tr>
<td>Mitochondria from mouse liver</td>
<td>$-0.30$ (several more $+50.0$)</td>
</tr>
<tr>
<td>Single protein layer on porous agarose beads</td>
<td>$-0.005$ to $-0.02$</td>
</tr>
<tr>
<td>Antigen-antibody double layers on Sepharose 4B beads</td>
<td>$+0.003$ to $+0.08$</td>
</tr>
<tr>
<td>Spinal cord liquid</td>
<td>$-2.63$ to $-2.80$</td>
</tr>
</tbody>
</table>

The accuracy of magnetic measurements relatively environmental fluid having susceptibility value $-0.72 \times 10^{-6}$ (CGS units) $= -4 \pi \times 0.72 \times 10^{-6}$ (SI units) reached $0.5 \times 10^{-11}$ (CGS units).
like the phase crystalline transitions could facilitate the penetration of light considered in works also by Karu \textsuperscript{18} through the tubes of a collagen acupuncture network. Polarization microscopy of reflected, transmitted, and scattered light combined with CCD recording of static and dynamic patterns, and image view analysis have given rise to a new method for the observation and measurements of magnetic separation processes in gradient magnetic fields, and of the magnetic alignment of biological microparticles, nucleoprotein granules, invertebrate eggs, and macromolecular aggregates of ferritin. A nucleoprotein granular suspension from the blood capillary system of a rabbit was studied in connections with the problem of the medical treatments of acupuncture points by polarized light, steady and gradient magnetic fields, and acoustic stimuli.\textsuperscript{17,21} The liquid crystalline arrangement, magnetic susceptibility changes, the appearance of optical and magnetic anisotropy in gradient magnetic fields, concentrating, light illumination, and other influences under time evolution were revealed.

The orientation behavior of tubuline,\textsuperscript{9} invertebrate eggs,\textsuperscript{5} retinal and colloidal core-shell rods,\textsuperscript{10} and DNA molecules,\textsuperscript{16} is becoming of particular interest. The relatively rigid DNA double helix with the base pairs stacked coplanar to each other and perpendicular to the long filament axis splits into two statistical, flexible, random coils, where the bases are free to rotate. In the denaturation phase transition region around the melting temperature, the magnetic orientation decreases to zero, and the observed decrease of optical and diamagnetic anisotropy led to the assumption that the base pairs are responsible for the magnetic alignment of DNA (Ref. 16). These base pairs in part consist of aromatic, i.e., diamagnetically anisotropic, rings and consequently tend to orient their planes parallel to field. The DNA coils in nucleosome supramolecular condensed chromatin structures tend to orient their planes parallel to \( H_0 \), which is in agreement with our observation.

The use of HGMS combined with polarization microscopy, CCD recording, and image analysis of arising the ordered crystalline structures enabled detection of small amounts of sample, and so localization of the visual field on the gradient magnetic areas. The magnetic field caused the alignment of nucleoprotein granules, their concentrating, phase transitions, and magnetic susceptibility sign changes in the forming structures. Mechanical stimulation combined with HGMS facilitated liquid crystalline transitions, and caused resonant mechanical vibrations of aggregates.\textsuperscript{21} During the oxidative and aging processes, the fragmentation of samples led to the formation of a liquid crystalline state with completely new magnetic properties, in particular, diamagnetic crystalline clusters were revealed in a gradient magnetic field. New methods were developed to detect isotropic-crystalline phase transitions while recording the orientation of DNA granule aggregate in a high-gradient magnetic field. In the first step (see Fig. 6), granules were detected as aggregates having a discoid form and revealing positive diamagnetic anisotropy, orienting themselves with their plane along the magnetic field in diamagnetic separation on the wire. The second orientation of aggregates changed their form to a cylindrical-rod-like form, then to conic forms, which were oriented perpendicular to the field and to the wire, thus manifesting negative diamagnetic anisotropy, which is the characteristic behavior, in our case, of nucleoprotein and DNA molecules. A new method was presented for the liquid phase state testing of a DNA-granule colloidal suspension of a nucleus containing microcells, whose aggregates are based on high-gradient magnetic separation and microscopic screening of magnetic accumulation with time evolution. Birefringence properties of the crystals that result from the DNA granules concentrated by high-gradient magnetic separation were studied with image analysis.

Our study demonstrated various applications of a magnetic separation analysis of bioparticles moving or collecting in a gradient magnetic field with the value \((H \times \text{grad}H)\) by three orders more than for technical systems without wires reported by other authors.\textsuperscript{22–28} who measured the red blood cell iron and hemoglobin through the oxidation of iron ions irreversibly into third-valence states. One group observed the migration of lymphocytes for several hours in a weak gradient, while we can measure the magnetic susceptibility of native cells for several seconds. The diamagnetic capture of live bioparticles in the gradient magnetic field in our method lasted 1 to 20 s, providing measured values such low as \(10^{-9}\) (SI units) and a standard deviation 5%. Other magnetophoretic reported data\textsuperscript{26} suggested the possibility to be applied for biotechnology using fluorescent immunomagnetic sorbent beads containing superparamagnetic cores for the collection of particles in a visual field. Our purpose was to understand the role of magnetic and optical changes that occur permanently and play the main role in biosensory mechanisms of living systems. Recent new studies have shown the existence of the acupuncture meridian threadlike structures in the human body.\textsuperscript{17,20} After light stimulation of the body, “light channels” appear to be identical to meridians in traditional Chinese medicine.

### 5 Conclusions

To measure the magnetic susceptibility and magnetic moments of small numbers of weakly magnetic particles, both diamagnetic and paramagnetic, we performed measurements in a high-gradient magnetic field (HGMF), which was realized with portable, inexpensive devices, mounted in the visual field of inverted or reflected microscopes. The study of magnetic materials in living systems is very appropriate for the development of new methods for magnetic diagnostics and magnetotherapy. Living systems contain the full interval of magnetic substances (Table 3), from ferromagnetic crystals
and superparamagnetic nanoparticles to diamagnetic anisotropic macromolecular structures.

Immunocomplex assays, which are very important in diagnostics and cancer immunology, are developed on magnetic separation principles. The suggested method could detect a protein antigen within several seconds. All these branches of investigation are the subject of great interest in the world.

Acknowledgments

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References