Enhanced photoacoustic stability of gold nanorods by silica matrix confinement

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Abstract. Photoacoustic tomography (PAT) has garnered much attention for its high contrast and excellent spatial resolution of perfused tissues. Gold nanorods (GNRs) have been employed to further enhance the imaging contrast of PAT. However, the photon fluences typically needed for PA wave induction often also result in GNR shape changes that significantly reduce the efficiency of acoustic wave generation. In this work, we propose, synthesize, and evaluate amorphous silica-coated gold nanorods (GNR-Si) in an effort to improve contrast agent stability and ameliorate efficiency loss during photoacoustic (PA) wave induction. TEM and optical absorption spectra measurements of GNR and GNR-Si show that encasing GNRs within amorphous silica provides substantial protection of nanorod conformation from thermal deformation. PA signals generated by GNR-Si demonstrate considerably greater resistance to degradation of signal intensity with repetitive pulsing than do uncoated GNRs, thereby enabling much longer, high-contrast imaging sessions than previously possible. The prolongation of high-contrast imaging, and biocompatibility and easy surface functionalization for targeting ligands afforded by amorphous silica, suggest GNR-Si to be potentially significant for the clinical translation of PAT. © 2010 Society of Photo-Optical Instrumentation Engineers. [DOI: 10.1117/1.3292574]

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Recently there has been considerable interest in employing the photoacoustic effect to tomographically image biological tissues, both in vitro and in vivo. Termed photoacoustic tomography (PAT), the emerging technique employs pulses of light to transiently generate thermoelastic tissue expansion, with a magnitude that depends on the local tissue’s optical and thermoacoustic properties. The excited tissue’s expansion and ensuing contraction give rise to a propagating broadband ultrasonic wave whose intensity is proportional to photon absorption and whose travel time is proportional to depth, subject to intervening tissue absorption, scattering, and dispersion. Measurements of the acoustic wave’s intensity and arrival time are made outside the sample with piezoelectric pressure transducers, from which 3-D maps of photon absorption can be reconstructed. To enable deeper tissue penetration of photons for acoustic wave production, near-infrared light of wavelengths between 700 and 900 nm is used to exploit the modest translucency of most mammalian tissues at these wave lengths.
wavelengths. Inherent imaging contrast arises from differences in the optical absorption of various tissues and their constituents, like the ten-fold difference in light absorption of blood from that of surrounding tissue.

Since the number of intrinsic imaging contrast agents found in vivo is limited, a variety of exogenous contrast agents have been developed to broaden the repertoire of PA targets that include near-infrared fluorescent dyes (e.g., indocyanine green), carbon nanotubes, and metallic nanoparticles. For this work, we selected gold nanorods (GNRs) as the platform of study, due to their easily tuned optical absorption and high biocompatibility. Previous studies have shown the longitudinal surface plasmon resonance of GNRs to be highly sensitive to nanorod size and aspect ratio, with strong correlation of those parameters to the wavelength and efficiency of photon absorption. Under pulsed laser excitation of sufficient energy to generate PA waves in vivo, however, transmission electron microscopy (TEM) studies have found significant permanent deformation of nanorods that correlated with adverse changes in the nanoparticle’s optical absorption, with increasing cumulative energy deposition resulting in blue-shifted, less intense absorption spectra for the longitudinal surface plasmon. Although several groups have exploited this laser-induced shift in the GNRs’ peak absorption of near-infrared light to characterize fluid flow, the consequential distortion and reduction of PA signal intensity severely limits its applicability to studies that permit replenishment of the contrast agent.

To ameliorate thermally induced deformation of GNRs used in PAT, we postulated that encasing the GNRs in a layer of amorphous silica (GNR-Si) might prove helpful; that, because of its rigidity and optical transparency at near-infrared wavelengths, silica might provide the GNR with the mechanical stability necessary to preclude deformation while minimizing photon absorption and scattering within itself. Moreover, amorphous silica’s biocompatibility and biodegradability are well established, with it having been successfully used in a number of bioimaging and biosensing applications.

In this work, bare GNRs with aspect ratios ≈4 were prepared as described by Nikoobakh and El-Sayed and then coated with porous silica using the single-step synthesis reported by Gorelikov and Matsuura. GNR and GNR-Si (approximately 10^{15} particles/L) morphologies were characterized with transmission electron microscopy (Hitachi H-7650 TEM), as shown in Fig. 1(a).

As shown in Fig. 1(b), a Nd:YAG laser (LS-2132U, LOTIS TII, Minsk, Belarus) operating at 1064 nm was used to pump a widely tunable Ti:sapphire laser (CF-125, SOLAR TII, Minsk, Republic of Belarus), delivering 16 to 18-nsec pulses of 800-nm light at a pulse repetition frequency of 10 Hz. The laser’s output beam was attenuated with a lens before entering a water tank that held a Rexolite 1422 (San Diego Plastics, California) tube phantom bearing 100 mL of GNR or GNR-Si in water. PA signals were detected with a commercial piezoelectric transducer (V303, Panametrics) immersed in the tank that possessed a center frequency of 1 MHz and a line focal depth of 17.8 mm. The pressure transducer’s output was amplified with an ultrasonic receiver (5077PR, Panametrics, Waltham, Massachusetts) and recorded at a rate of 10 Msamples/sec on a PC-mounted, laser-triggered data acquisition card (CompuScope 14200, Gage, Lachine, Quebec, Canada).

In Fig. 2 we present the PA signal intensity of GNR and GNR-Si as a function of repeated laser excitation for two different photon fluences. Signals were normalized to their respective initial values (i.e., maxima) and fit to the single exponential curves. Early (<1 sec) PA signal intensities of GNR and GNR-Si were found to be approximately equal to one another [PA signal intensities were 48.81 for GNR and 44.28 for GNR-Si in Fig. 2a] and 123 for GNR and 126.3 for GNR-Si in Fig. 2(b), suggesting that the thermal-acoustic impedance mismatch—introduced by the silica layer between the metal nanorod and its aqueous surroundings—was of little consequence to the formation of acoustic waves. When exposed to 21.20 mJ/cm^2 of coherent 800-nm light [Fig. 2(a)], the PA signal intensity of GNR decays rapidly during the first 30 sec to 57% of its initial value, and decreases further to 45% of its initial value after 2 min. When GNR-Si is subjected to identical laser irradiation, however, the PA signal decays only to 86% of its initial maximal intensity during the first 30 sec, and remains relatively constant at that level thereafter.

Increasing the excitation photon fluence further resulted in considerable degradation of the PA response from both GNR and GNR-Si [Fig. 2(b)]. At exposures of 31.7 mJ/cm^2 [the ANSI Z136 maximum permissible exposure (MPE) of human skin to visible/near-visible coherent light] the PA signal intensity of GNR decays very rapidly to 45% of its initial value within the first 10 sec, and decreases further to 37% of its initial strength after 30 sec. By comparison, 31.7-mJ/cm^2 exposures of GNR-Si resulted in the average PA signal intensity decaying to 58% of its initial value after 10 sec, and remaining stable thereafter—quite unlike the continuing decay exhibited by GNR’s PA signal. Increasing the photon fluence to 103.27 mJ/cm^2 showed no effect on the PA signal generated by GNR-Si, with an intensity approximately 60% of that of its initial value after 30 sec, and remaining stable thereafter.

To verify that these changes in PA signal intensity with time and photon fluence arose from the deformation of nanorods and not some other source, we made a series of TEM and optical absorption spectrum (DU800 UV, Beckman) measurements, as shown in Fig. 3. When subjected to exposures of 21.2 mJ/cm^2, TEM studies revealed that a significant fraction of the original bare GNRs [upper left, Fig. 3(a)] had deformed into rounder, lower aspect ratio entities, while spectroscopy showed the longitudinal surface plasmon’s absorption peak had flattened and blue-shifted [center left, Fig. 3(a)]. GNR-Si [upper right, Fig. 3(b)], however, did not exhibit such behavior at these levels of irradiation, being able to, more or less, maintain their original morphology and absorption spectrum [center right, Fig. 3(b)].

When GNRs were subjected to photon fluences of 31.12 mJ/cm^2, TEM studies showed that the overwhelming majority of bare GNR had distorted into either oblate/prolate spheroids or bent, v-shaped cylinders—with the longitudinal plasmon’s absorption peak height reduced to roughly that of the transverse plasmon’s 520-nm (aspect ratio-independent) peak [left, Fig. 3(b)]. In contrast, TEM morphology and absorption spectra measurements of GNR-Si demonstrated rela-
tively modest differences between 21.20 and 31.12 mJ/cm² irradiated specimens, with less than a 10% height reduction and a ~60-nm blue-shift in the longitudinal plasmon’s absorption peak for higher doses [right, Fig. 3b]. Thus, postirradiated GNR-Si retained many of the structural and optical attributes of preirradiated GNR-Si, consistent with PA signal intensity measurements.

Unfortunately, both sample preparation and practical limits on the number of samples/images that can be analyzed often result in TEM yielding statistically poor representations of nanorod aspect ratio distributions. To address this shortcoming in our analysis based on TEM measurements, we fitted our empirically derived inhomogeneously broadened optical absorption spectra to multiple calculated homogeneously broadened absorption spectra that were each synthesized via Mie-Gans theory for a collection of single-aspect-ratio nanorods. The weighting/frequency of each aspect ratio’s calculated spectra used in fitting measured spectra was determined iteratively by a least squares solution of minimum norm between discretized forms of the measured and calculated spectra at a given wavelength.

Fig. 1 (a) TEM images and Mie-Gans theory spectra fits for GNRs [left, 1(a)] and GNR-Si [right, 1(a)] prior to laser irradiation. (b) Photoacoustic imaging system used in the studies.
The individually weighted, calculated absorption spectra and their summation for different laser exposures of GNR and GNR-Si are shown in Figs. 1 and 2. As the graphs illustrate, the percentage of GNRs with aspect ratios between 3.75 and 4.75, which reflect nanorods with longitudinal surface plasmon absorption peaks at wavelengths close to that of our GNR before laser deformation (800 nm), decreased from 61.0 to 24.3% after 21.20-mJ/cm² exposure, and down to 17.5% after 31.12-mJ/cm² exposure, while corresponding exposures of GNR-Si resulted in decreases from 64.3 to 46.4 to 35.8%, respectively. We note, however, that care must be exercised when interpreting the aspect ratio distributions of more thermally deformed nanorods via this method since, as TEM images like that of Figs. 3a and 3b show, a number of nanorods become angular/bent with repeated laser irradiation, and thus no longer satisfy the smoothness of topology assumed by either the Mie theory or Gans extension of the Mie theory.

Indeed, under such circumstances, rigorous application of the discrete dipole approximation (DDA) may be more appropriate than Mie-Gans modeling, albeit at far greater computational expense, since both scattering and absorption efficiencies must be considered (due to their comparable magnitudes), and large numbers (10⁴ < N < 10⁶) of point dipoles [perhaps on nonperiodic lattices that preclude fast Fourier transform (FFT) simplification, or inclusion of higher-order multiple excitations] will be needed to approximate light scattering/absorption by irregular surfaces (e.g., vertices, crevices) and ensure convergence of optical properties. Moreover, surface plasmons and intraband transitions tend to
obscure many of the spectral fine structure that otherwise characterize light scattering/absorption by angular metallic nanoparticles. Thus there is considerable a-priori concern as to the ability of DDA methods to uniquely describe the distribution of both gold nanoparticle sizes and shapes responsible for a particular optical signature.

To determine if the stability conferred by silica encasement on nanorod morphology actually resulted in improved PA imaging contrast and stability, the system shown in Fig. [1] was used but with the ultrasonic transducer, now driven by a precision linear stage (HR8, Nanomotion, Yokneam, Israel). A phantom comprised of a 0.5-cm-diam tube filled with cotton and either GNR or GNR-Si was placed within the tank, orthogonal to the transducer. During a scan, discrete PA measurements were made as the transducer was stepped in 0.1-mm increments across the tube’s minor axis to provide 1-D cross sectional views. PA contrast assessment was made by including cotton packing material within the phantom to provide a source of nonuniform GNR or GNR-Si distribution. For these nontomographic studies, the MPE photon fluence of 31.12 mJ/cm² was used to provide clinically relevant signal-to-noise ratios (SNRs) in imaging.

In Fig. 4(a) we show PA intensity images of tubes filled with GNR (top) and GNR-Si (bottom), derived from 18-sec samplings that began t=0 sec (left) and t=108 sec (right) after commencement of continuous laser irradiation. As the circled regions of Fig. 4(a) emphasize, GNR suffered considerable loss of image detail as laser exposure continued—a situation not observed in GNR-Si, whose PA image appears largely unchanged at later time points. Sequential comparison of the intensity in each image of a given tube also revealed a similar loss of information with time, with the intensity of GNR decreasing by 3 dB after 90 sec, while that of GNRs at Si decreased by only 0.7 dB from their respective initial values [Fig. 4(b)], consistent with earlier, nonimaging PA signal intensity measurements.

In summary, we have synthesized and evaluated both bare and amorphous silica-coated GNRs in an effort to determine whether the confinement of GNRs by rigid silica would improve their utility as photoacoustic contrast agents. TEM and optical absorption studies of bare GNR versus GNR-Si revealed the latter to have significantly greater resistance than bare GNR to thermal deformation under repeated laser irradiation, without diminution of the longitudinal surface plasmon’s absorption efficiency, even at photon fluences more than three times the MPE. PA signals generated by GNR-Si similarly demonstrated commensurately greater resistance to degradation of signal intensity with repetitive pulsing than did uncoated GNRs, enabling much longer, high-contrast imaging sessions than previously possible. The prolongation of high-contrast imaging and easy surface functionalization for targeting ligand binding afforded by the amorphous silica coating suggest GNR-Si to be a significant step toward the clinical translation of PAT.

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