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Abstract. Blast-induced traumatic brain injury has been a recent major concern in neurotraumatology. However, its pathophysiology and mechanism are not understood partly due to insufficient information on the brain pathophysiology during/immediately after shock wave exposure. We transcranially applied a laser-induced shock wave (LISW, ~19 Pa · s) to the left frontal region in a rat and performed multispectral imaging of the ipsilateral cortex through a cranial window (n = 4). For the spectral data obtained, we conducted multiple regression analysis aided by Monte Carlo simulation to evaluate vascular diameters, regional hemoglobin concentration (rCHb), tissue oxygen saturation (StO2), oxygen extraction fraction, and light-scattering signals as a signature of cortical spreading depolarization (CSD). Immediately after LISW exposure, rCHb and StO2 were significantly decreased with distinct venular constriction. CSD was then generated and was accompanied by distinct hyperemia/hyperoxemia. This was followed by oligemia with arteriolar constriction, but it soon recovered (within ~20 min). However, severe hypoxemia was persistently observed during the post-CSD period (~1 h). These observations indicate that inadequate oxygen supply and/or excessive oxygen consumption continued even after blood supply was restored in the cortex. Such a hypoxic state and/or a hypermetabolic state might be associated with brain damage caused by a shock wave. © The Authors. Published by SPIE under a Creative Commons Attribution 4.0 Unported License. Distribution or reproduction of this work in whole or in part requires full attribution of the original publication, including its DOI. [DOI: 10.1117/1.JBO.24.3.035005]

Keywords: blast-induced mild traumatic brain injury; multispectral imaging; diffuse reflectance; cortical spreading depolarization; hemodynamics; light scattering.

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1 Introduction

In recent years, the risk of military personnel and civilians suffering from blast-induced traumatic brain injury (bTBI) has been increasing due to frequent attacks using improvised explosive devices.1,2 The majority of the patients (>90%) have been categorized as having mild to moderate TBI; most of them do not show any abnormality by conventional structural brain imaging, such as computed tomography. However, many of the patients suffer from higher-order brain dysfunction as well as other chronic symptoms, such as headaches, photophobia, and insomnia, which considerably decrease their quality of life.1,4 Despite extensive studies, the pathophysiology and mechanisms of such bTBI have not been fully elucidated. This is largely due to, in addition to the complex interactions of the brain with a shock wave, the lack of information on physiological, structural, and functional changes in the brain during and immediately after shock wave exposure in vivo. Since initial events caused by a shock wave can trigger cascading processes leading to pathological consequences in the chronic phase, real-time observation of the cerebral responses to a shock wave is crucial for understanding the mechanism of bTBI. However, real-time monitoring and assessment of the brain exposed to a shock wave is generally difficult, since sensing or detecting systems can be considerably affected by an experimental blast or shock wave. Here, optical technologies can play important roles.

We have been using a laser-induced shock wave (LISW) to investigate the pathophysiology and mechanism of bTBI.5,6 The advantages of the use of an LISW for bTBI research include its safety, ease in handling, and ability to precisely control the shock wave energy and size. In addition, the confined nature of the shock wave energy of an LISW enables site-specific or targeted application to animals, which is useful because blast injuries resulting from systemic blast exposure are often too complex to analyze or understand the mechanism of injuries.11,13 Furthermore, sensors can be safely placed in the vicinity of the site of LISW exposure, enabling real-time monitoring of the brain pathophysiology. In our previous study, we performed transcranial real-time diagnosis of the rat brain exposed to an LISW by using optical fiber-based diffuse reflectance spectroscopy (DRS).5,14 We observed that, immediately after LISW application, cortical spreading depolarization (CSD) was generated at the site of LISW application and propagated over the entire ipsilateral cortex. This was accompanied by distinct hyperemia/hyperoxemia, but after CSD propagation, long-lasting hypoxemia/oligemia was observed for up to ~3 h. CSD is a wave of near-complete neuronal and glial
depolarization associated with massive transmembrane ionic and water movements, therefore imposing a considerable metabolic workload on the tissue for restoration of imbalanced ions. Thus, CSD has received much attention to understand the mechanism and symptoms of various brain diseases and injury.\(^{13-17}\) To the authors’ best knowledge, however, there has been no report on the occurrence of CSD in bTBI animal models.

Importantly, it has been reported that the consequences of CSD vary depending on the baseline metabolic and hemodynamic states in the tissue. In addition, since CSD is a neurovascular event, bidirectional interactions between the parenchymal tissue and cerebral vasculature during and after CSD are crucial for understanding the pathological consequences of CSD in its disease.\(^{5,6}\) However, detailed characteristics of shock wave-induced CSD are not elucidated. In this study, we performed multispectral imaging of the cortex of the rat brain that was transcranially exposed to an LISW and quantified two-dimensional cerebral hemodynamics and tissue scattering parameters using multiple regression analysis aided by Monte Carlo simulation for the spectral data obtained.\(^{18,19}\) We assessed morphological changes of vessels and visualized distributions of regional hemoglobin concentration (\(c_{\text{Hb}}\)) and tissue oxygen saturation (\(S_{\text{O2}}\)) for three time regimes: (i) just after LISW application (before the start of CSD), (ii) during CSD propagation, and (iii) after CSD propagation. Importantly, vascular type-dependent behaviors were observed, and the results were therefore analyzed for each vascular type (arteriolar, venular, and capillaries). Oxygen extraction fraction (OEF) was also evaluated as a clinically relevant parameter related to cerebral oxygen metabolism. To the authors’ best knowledge, there is no report on such detailed spatiotemporal responses of the brain exposed to a shock wave in such super acute phases. On the basis of all the results obtained, we discuss possible mechanisms of brain tissue alteration after exposure to a shock wave.

## 2 Materials and Methods

All requests for animals and intended procedures in the present study were approved by the Ethics Committee of Animal Care and Experimentation, National Defense Medical College, Japan (permission numbers: 13024, 16036). All animal experiments were performed under anesthesia, as described later, and all efforts were made to minimize suffering.

### 2.1 Generation and Characteristics of a Laser-Induced Shock Wave

Figure 1(a) shows the method for generation of an LISW. A laser target, which was a light-absorbing material (0.5-mm-thick natural black rubber disk) on which an optically transparent material (1.0-mm-thick polyethylene terephthalate sheet) was adhered with adhesive, was placed on the tissue. The laser target was irradiated with a high-intensity laser pulse from a 532-nm Q-switched Nd:YAG laser [pulse width, 6-ns full-width at half-maximum (FWHM); Brilliant b, Quantel, Les Ulis Cedex, France]. The laser pulse was absorbed by the black rubber to induce plasma and its expansion generated a shock wave (LISW). We previously observed that an LISW generated at a laser fluence of \(1.25\) \(\mu\)J/cm\(^2\) with a spot diameter of \(4\) mm (peak pressure: \(\sim 104\) MPa, impulse: \(\sim 19\) Pa \(\cdot\) s, positive pulse duration: \(\sim 0.8\) \(\mu\)s) caused CSD in all of the five rats examined.\(^{5}\)

In this condition, transcranial application of a single LISW to a rat brain caused either no or limited hemorrhage in the cortical surface. On the basis of the histopathological analysis, we also confirmed that there was neither thermal damage nor mechanical damage (contusion) in the cortex at the same condition (data not shown). Thus, this condition can be regarded as mild or mild-to-moderate injury and was used in the present study.

The peak pressure of an LISW is 2 to 3 orders of magnitude higher than that of medically relevant actual explosion shock waves, while the positive pressure duration of an LISW is 2 to 3 orders of magnitude shorter. As a result, the impulse (the time-integrated positive pressure component), which is known to be one of the most important parameters to determine tissue damage, is in the same order as that of relevant actual explosion shock waves.\(^{20}\) Thus, we assume that the condition we used is valid to investigate the brain pathophysiology associated with bTBI in rats.

### 2.2 Surgical Preparation and LISW Application

Male Sprague–Dawley rats weighing 300 to 400 g (10 to 12 weeks of age, \(n = 4\), Japan SLC, Inc., Shizuoka, Japan) were anesthetized with an intraperitoneal injection of pentobarbital sodium (50 mg/kg bolus and \(\sim 15\) mg/kg/h maintenance)
and placed in a stereotactic frame. After shaving the head, the scalp was incised at the midline and the parietal bone was exposed. During experiments, adequacy of anesthesia was checked by the absence of a response to tail pinch. For multispectral imaging, a cranial window with a diameter of ~4.5 mm was prepared with the dura intact in the left parietal bone with a trephine (Roboz Surgical Instrument, Co., Inc., Gaithersburg, Maryland) under a surgical microscope (OPMI Pico/S100, Carl Zeiss, Germany). A disk of transparent polymethyl methacrylate with a diameter of 4.4 mm and a thickness of ~0.5 mm was inserted within the window and sealed with dental cement. Figure 1(b) shows a schematic illustrating positions of the cranial window and the site of LISW application. The distance between the center of the LISW and the center of the cranial window was ~7 mm. After window preparation, the rat was moved to the position for imaging and placed in another stereotactic frame. The body temperature was kept constant at 37.0°C ± 0.5°C by using a temperature-controlled body mat, and arterial oxygen saturation (SpO₂) was continuously monitored at the hindlimb with a pulse oximeter (8600V, Nonin Medical, Inc., Plymouth, Minnesota).

2.3 Multispectral Imaging

Through the window described above, we performed multispectral imaging of the cortex to observe CSD, morphology of vessels (changes in diameters of cortical surface vessels), rCHb, and StO₂ before, during, and after LISW application. Figure 1(c) shows a diagram of the experimental setup. The parietal imaging region was illuminated with the light from a halogen lamp (HL 100E, HOYA-Schott, Tokyo, Japan); the wavelength of the light was changed sequentially using a rotating filter wheel (FW102C, Thorlabs, Newton, New Jersey) with six bandpass filters. The center wavelengths of the filters were 520, 540, 560, 580, 600, and 800 nm and their bandwidths were 10 nm FWHM (520 to 600 nm) or 40 nm FWHM (800 nm); the diameters of the filters were ~25 mm (Thorlabs). Diffusely reflected light from the cortex was captured with a CCD (12 bit, 1034 × 779 pixels, 1/30 s/frame, scA1000-30fm, Basler AG, Ahrensburg, Germany) equipped with a unifocal lens (f = 50 mm, M118FM50, Tamron Co., Ltd., Saitama, Japan). The filter wheel was rotated every 1.5 s using accessory software (Thorlabs), and image acquisitions were synchronized with the filter rotations using an output signal from the filter wheel and LabVIEW software (National Instruments, Austin, Texas), the image acquisition rate being 1 image per 1.5 s at each wavelength. Since images at 5 of the 6 wavelengths were used for quantitative analyses as described later, the actual imaging rate was 1 image per 7.5 s. A laser target was placed on the skull in the frontal region and held with forceps; ultrasound gel was applied between the bottom of the target (rubber) and the skull surface for acoustic impedance matching. Baseline images were acquired at least for 2 min, and the target was then irradiated with a single laser pulse to produce an LISW, and imaging was continued for ~60 min. After imaging, a certified diffuse reflectance standard plate (SRS-99-020, Labsphere, Inc., North Sutton, New Hampshire) was placed at the same position and its image was acquired under the same imaging conditions. All multispectral images acquired for the cortex were divided by images acquired for the reflectance standard at the corresponding wavelengths, which were used as diffuse reflectance r(λ) images to calculate optical and physiological parameters, as described below.

2.4 Estimation of Cerebral Absorption (Hemodynamics) and Scattering Parameters

We performed multiple regression analysis with Monte Carlo simulation to map the reduced scattering coefficient, which was used as an optical marker for CSD, and absorption coefficient as well as rCHb and StO₂, using MATLAB software (Mathworks, Natick, Massachusetts). Images at 520, 540, 560, 580, and 600 nm were used for assessing the reduced scattering coefficient and concentrations of oxygenated hemoglobin (C_HbO) and deoxygenated hemoglobin (C_HbR), while images at 800 nm were used to determine the propagation velocity of CSD, as described below. The method is basically the same as that used in our previous study. Figure 2 shows the flow diagram of the method. Absorbance spectrum A(λ) was calculated from diffuse reflectance r(λ) based on the following equation:

\[ A(\lambda) = -\log_{10} r(\lambda), \]

where \( r(\lambda) \) corresponds to the diffuse reflectance spectrum normalized by the incident light spectrum. Because attenuation due to light scattering can be treated as pseudochromophore, the absorbance spectrum A(λ) can be approximated as the sum of attenuations due to absorption and scattering in the brain as follows:

\[ A(\lambda) = C_{HbO} A(\lambda, C_{HbO}, C_{HbR}, \mu'_s) + C_{HbR} A(\lambda, C_{HbO}, C_{HbR}, \mu'_s) + D(\lambda, \mu'), \]

where C is the concentration, \( l \) is the mean path length, \( \epsilon \) is the molar extinction coefficient, and \( D(\lambda, \mu') \) indicates attenuation due to light scattering in the brain. The subscripts HbO and HbR denote oxygenated hemoglobin and deoxygenated hemoglobin, respectively. The absorption coefficient of the cortical tissue was assumed to depend only on \( C_{HbO} \) and \( C_{HbR} \) as follows:

\[ \mu_a(\lambda) = C \epsilon(\lambda) = C_{HbO} \epsilon_{HbO}(\lambda) + C_{HbR} \epsilon_{HbR}(\lambda). \]

The total hemoglobin concentration \( C_{HbT} \), also referred to as \( CHb \), is defined as the sum of \( C_{HbO} \) and \( C_{HbR} \) as follows:

\[ C_{HbT} = C_{HbO} + C_{HbR}. \]

The tissue oxygen saturation is determined as follows:

\[ StO_2 = \frac{C_{HbO}}{C_{HbO} + C_{HbR}} \times 100. \]

The reduced scattering coefficient of the brain tissue was assumed to be in the following form:

\[ \mu'_s(\lambda) = a \lambda^{-b}, \]

where the coefficient a represents the scattering amplitude and b is the scattering power being constant at 1.38. By using \( A(\lambda) \) as a response variable and known extinction coefficients of oxygenated hemoglobin \( \epsilon_{HbO}(\lambda) \) and deoxygenated hemoglobin \( \epsilon_{HbR}(\lambda) \) as predictor variables, the multiple regression analysis based on the modified Lambert–Beer law (MRA1) can be applied to Eq. (2) as follows:
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\( A(\lambda) = a_{\text{HBO}}\bar{e}_{\text{HBO}}(\lambda) + a_{\text{HBR}}\bar{e}_{\text{HBR}}(\lambda) + a_0, \)

where \( a_{\text{HBO}}, a_{\text{HBR}}, \) and \( a_0 \) are the regression coefficients. The regression coefficients \( a_{\text{HBO}} \) and \( a_{\text{HBR}} \) describe the degree of contribution of each extinction coefficient to \( A(\lambda) \) and are closely related to the concentrations \( C_{\text{HBO}} \) and \( C_{\text{HBR}} \), respectively. The regression coefficient \( a_0 \) is expressed as follows:

\[ a_0(\lambda) = \bar{A} - \bar{e}_{\text{HBO}}a_{\text{HBO}} - \bar{e}_{\text{HBR}}a_{\text{HBR}}, \]

where \( \bar{A}, \bar{e}_{\text{HBO}}, \) and \( \bar{e}_{\text{HBR}} \) are the averages of \( A, e_{\text{HBO}}, \) and \( e_{\text{HBR}} \) over the wavelength range, respectively. \( a_0 \) represents the bias component of \( A(\lambda) \). Thus, \( a_0 \) describe the degree of contribution of the attenuation due to light scattering in the brain to the absorbance spectrum \( A(\lambda) \) and hence, it is related to the coefficient \( a \) in Eq. (6). At the same time, \( a_0 \) is also affected by absorption coefficient of the brain, since \( A(\lambda) \) is generally a function of the tissue absorption coefficient and reduced scattering coefficient.

As described above, the regression coefficients \( a_{\text{HBO}}, a_{\text{HBR}}, \) and \( a_0 \) are related to the concentration of oxygenated hemoglobin \( C_{\text{HBO}} \), that of deoxygenated hemoglobin \( C_{\text{HBR}}, \) and the coefficient \( a \), respectively. However, \( a_{\text{HBO}}, a_{\text{HBR}}, \) and \( a_0 \) are not determined by a unique regression coefficient when using only MRA1. Thus, we used further multiple regression analysis to estimate the values of \( C_{\text{HBO}}, C_{\text{HBR}}, \) and \( a \) on the basis of the combination of regression coefficients \( a_{\text{HBO}}, a_{\text{HBR}}, \) and \( a_0 \) that were obtained from MRA1. In this analysis, \( C_{\text{HBO}}, C_{\text{HBR}}, \) and \( a \) were regarded as response variables, and the three regression coefficients \( a_{\text{HBO}}, a_{\text{HBR}}, \) and \( a_0 \) in Eq. (7) were regarded as predictor variables. Their relations are written as \( C_{\text{HBO}} = \beta_{\text{HBO}}a_{\text{HBO}} + \beta_{\text{HBR}}a_{\text{HBR}} + \beta_a a_0 \), where \( a = \beta_a a_0 \), \( \alpha = [1, a_{\text{HBO}}, a_{\text{HBR}}] \), \( \beta_{\text{HBO}} = [\beta_{\text{HBO}}(0), \beta_{\text{HBO}}(1), \beta_{\text{HBO}}(2), \beta_{\text{HBO}}(3)] \), \( \beta_{\text{HBR}} = [\beta_{\text{HBR}}(0), \beta_{\text{HBR}}(1), \beta_{\text{HBR}}(2), \beta_{\text{HBR}}(3)] \), and \( \beta_a = [\beta_a(0), \beta_a(1), \beta_a(2), \beta_a(3)] \).

The symbol \(^T\) represents the transposition of a vector. The coefficients \( \beta_{\text{HBO}}, \beta_{\text{HBR}}, \) and \( \beta_a(i = 0, 1, 2, 3) \) are unknown and must be determined before analysis. We adopted MCS as the foundation to establish reliable values of \( \beta_{\text{HBO}}, \beta_{\text{HBR}}, \) and \( \beta_a \).
and $\beta_{\text{a};i}$. We used the MCS source code developed by Wang et al.,27 in which the Henyey–Greenstein phase function is applied to the sampling of the scattering angle of photons. The simulation model consisted of a single layer representing cortical tissue. In a single simulation of diffuse reflectance at each wavelength, 5,000,000 photons were randomly launched. The absorption coefficients of oxygenated hemoglobin $\mu_{\text{aHbO}}$ and deoxygenated hemoglobin $\mu_{\text{aHbR}}$ were obtained from the values of $\epsilon_{\text{HbO}}(\lambda)$ and $\epsilon_{\text{HbR}}(\lambda)$ in the literature,26 where the hemoglobin concentration of blood with a 44% hematocrit is 2326 $\mu$M of hemoglobin. The absorption coefficient $\mu_{\text{a}}$ converted from the concentrations $C_{\text{HbO}}$ and $C_{\text{HbR}}$ and the reduced scattering coefficient $\mu'_{\text{s}}$ specified by the coefficient $a$ were provided as inputs to the simulation, while the diffuse reflectance was produced as output. The input concentrations and coefficient $a$ and the output reflectance are helpful as the dataset in specifying the values of $\beta_{\text{a};i}$, $\beta_{\text{aHbO}}$, and $\beta_{\text{aHbR}}$ statistically for determining the absolute values of $C_{\text{HbO}}$, $C_{\text{HbR}}$, and $a$. The five different values of 40,172, 60,258, 80,430, and 120,516 were calculated by multiplying the typical value26 of $a$ by 0.5, 0.75, 1.0, 1.25, and 1.5, respectively, and the reduced scattering coefficients $\mu'_{\text{s}}(\lambda)$ of the cortical tissue with the five different values were derived from the relation of Eq. (6). The sum of the absorption coefficients of HbO and HbR, $\mu_{\text{aHbO}} + \mu_{\text{aHbR}}(\mu_{\text{aHbT}})$, for $C_{\text{HbT}} = 4.65, 23.3, 116, 233$, and 465 $\mu$M was used as input to the cortical tissue in the simulation. Tissue oxygen saturation of $\text{SiO}_2$ was determined by $\mu_{\text{aHbO}}(\mu_{\text{aHbT}})$ and values of 0, 20, 40, 60, 80, and 100% were used for simulation. In total, 150 diffuse reflectance spectra at $\lambda = 520, 540, 560, 580$, and 600 nm were simulated under the various combinations of $C_{\text{HbO}}$, $C_{\text{HbR}}$, and $a$. MRA1 for the measured spectrum based on Eq. (7) generated the 150 sets of coefficient $a$, and $C_{\text{HbO}}$, $C_{\text{HbR}}$, and $a$. The coefficient vectors, also referred to as conversion vector, $\beta_{\text{HbO}}$, $\beta_{\text{HbR}}$, and $\beta_a$ were determined statistically by performing MRA2. Once $\beta_{\text{HbO}}$, $\beta_{\text{HbR}}$, and $\beta_a$ were obtained, $C_{\text{HbO}}$, $C_{\text{HbR}}$, and $a$ were calculated from $\alpha_{\text{HbO}}$, $\alpha_{\text{HbR}}$, and $\alpha_{\text{a}}$, which were derived from MRA1 for the measured reflectance spectrum, without the Monte Carlo simulation. Therefore, the spectrum of absorption coefficient $\mu_{\text{a}}(\lambda)$ and that of reduced scattering coefficient $\mu'_{\text{s}}(\lambda)$ was reconstructed by Eqs. (3) and (6), respectively, from the measured reflectance spectrum. The principle of the method was applied to the example of tissue phantoms;25 the correlation between the given value of $\mu_{\text{a}}$ and estimated value of $\mu_{\text{a}}$, and that between the given value of $\mu_{\text{s}}$ and estimated value of $\mu_{\text{s}}$ were shown. We also compared the results of calculation of changes in $C_{\text{HbO}}$ and $C_{\text{HbR}}$ with those obtained by the conventional modified Lambert–Beer law analysis (see Sec. 5). As described above, since the light scattering signal can be used as an optical marker of CSD, we used maps of scattering amplitude $a$, which is proportional to $\mu'_{\text{s}}$, to visualize propagation of CSD in this study. To determine the propagation velocity of CSD, reflectance images at 800 nm were used since they are also sensitive to light scattering changes due to CSD.3

2.5 Evaluation of Cortical Vessel Diameters

To assess changes in diameters of cortical arterioles and venules, we used reflectance images at 520 nm, which is an isosbestic point of HbO and HbR. For the cortical image, we placed a linear line that was perpendicularly across the centerline of the vessel of interest on LabVIEW software (see white and red lines on the top image in Fig. 5), creating an image intensity profile of the vessel. Vessel diameter was defined as the FWHM value of the profile.29,30,31 Arterioles and venules were distinguished on the basis of their diameters, positions, orientations, and values of $\text{SiO}_2$ before LISW application. In addition, we carefully distinguished cortical surface arterioles from the vessels in the dura matter on the basis of their diameters; the diameter of cortical surface arterioles is about two to three times larger than that of dural vessels ($\sim25 \mu$m) in their proximal portion in rats.23

2.6 Statistical Analysis

Statistical methods for data obtained in this study were selected on the basis of analysis of data normality and uniformity of variance using the Eksusu–Toukei 2012 statistical software package (Social Survey Research Information Co., Ltd., Tokyo, Japan). With regard to changes in diameters of pial arterioles and venules, results were compared using the nonparametric Dunn’s multiple comparisons test. Changes in $rC_{\text{Hb}}$ and $\text{SiO}_2$ were tested by the parametric Dunnett’s multiple comparisons test or the nonparametric Dunn’s multiple comparisons test.
Changes in OEF were tested by the parametric Tukey’s multiple comparisons test. The multiple comparison tests were performed using GraphPad Prism software (Prism 6 for Windows, GraphPad software, Inc., La Jolla, California); a difference was considered statistically significant when \( p < 0.05 \). Data are presented as means ± standard deviation (SD) when the parametric method was used, and data are presented as medians with ranges (in box plots) when the nonparametric method was selected.

### 3 Results

#### 3.1 Spatiotemporal Changes in Scattering Amplitude Associated with Cortical Spreading Depolarization

In our previous studies, we showed that light scattering signals can be used as an optical marker of CSD. Figure 3 shows...
difference images of scattering amplitude in the rat cortex after exposure to an LISW, where the image before LISW application is subtracted from the image at each timepoint (rat no. 3). After LISW application, a large venule showed an increase in scattering amplitude, and the region with increased scattering amplitude appeared in the parenchyma and in the pial arterioles in the right edge of the region of interest (RoI) (anterior side of the cortex) at 1 min 21 s after LISW application. The region with increased scattering amplitude in the parenchyma and arterioles was propagating like a wave to the left (posterior side), indicating a propagation of CSD. Such CSD-related scattering change was observed in all four rats examined, and only in one...

**Fig. 5** Time courses of changes in diameters of pial arterioles and venules after LISW application for rat no. 3; diameters were normalized by average diameters before LISW application. (a) Arteriole A anterograde to the direction of CSD propagation (A-1 to 3). (b) Arteriole B perpendicular to the direction of CSD propagation (B-1 to 3). (c) Venule C with a small diameter (C-1 to 3). (d) Venule D with a large diameter (D-1 to 3). Each position to measure the diameter is shown in the top image. Each graph on the left side is a magnified view of the bold-lined square of the graph on the right side. Arrowheads from \( t_1 \) to \( t_5 \) show timepoints for analysis in Figs. 6, 9, and 10.
rat (rat no. 4), the second consecutive wave of CSD was seen at 5 min 45 s after the start of the first wave. The propagation velocity of the wavefront of the first CSD wave in the parenchyma was $2.96 \pm 0.55$ mm/min (mean $\pm$ SD) ($n = 4$), being consistent with the results of our previous study and other studies using KCl injection or pinprick to induce CSD.

### 3.2 Vasculature Morphological Changes

Figure 4 shows images at 520 nm, an isosbestic wavelength of HbO and HbR, of pial arterioles and venules before and after exposure to an LISW in the same rat, as that shown in Fig. 3 (rat no. 3). Vessel type-dependent morphological changes were seen. A pial arteriole marked “A” that is running in the anterograde direction to the CSD propagation direction showed almost no change in diameter just after LISW application but showed distinct vasodilatation during the CSD propagation and then recovered soon after the CSD propagation [Fig. 4(a)]. Such vasodilatation during CSD was not observed in the pial arteriole marked “B” that is running perpendicular to the CSD propagation direction [Fig. 4(b)]. Pial venules observed through the cranial window were all anatomically perpendicular to or retrograde to the CSD propagation direction, while their initial diameters were greatly different. Thus, we divided them into two groups: smaller-diameter (less than 80 μm) pial venules such as that marked “C” and larger-diameter (80 μm or more) pial venules

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**Fig. 6** Relative changes in diameters of pial arterioles and venules at five time points after LISW application from the baseline for four rats examined; description for each time point is shown at the bottom of the graphs. (a) Arteriole A anterograde to the direction of CSD propagation. (b) Arteriole B perpendicular to the direction of CSD propagation. (c) Venule C with a diameter less than 80 μm. (d) Venule D with a diameter of 80 μm or more. The horizontal dashed line in each graph shows the baseline before LISW application. Numbers of positions for vessel diameter measurements in rat no. 1, no. 2, no. 3, and no. 4 are as follows: (a) arteriole A, 3, 6, 3, and 2; (b) arteriole B, 2, 0, 3, and 3; (c) venules C, 5, 4, 4, and 4; (d) venule D, 4, 2, 4, and 2, respectively. Data are presented as medians with ranges. Asterisks $(p < 0.05, \ast \ast p < 0.01)$ indicate a statistically significant difference (Dunn’s multiple comparisons test).
such as that marked “D.” Both types of venules showed distinct constriction immediately after LISW application regardless of their orientations and distances from the site of LISW application [Figs. 4(c) and 4(d)]; smaller-diameter venule C appeared to be more prominently constricted. The constrictions in both types of venules tended to recover during CSD propagation, but both types of venules reconstricted after the CSD propagation.

Figure 5 shows time courses of changes in diameters of each type of arteriole and venule for ∼60 min after LISW application (rat no. 3), for which we randomly selected three positions on each arteriole/venule shown in Fig. 4. The diameter of arteriole A slightly increased and then decreased by 10% to 15% over a period of 1 min after LISW application and then suddenly increased by a maximum of ∼80% concomitantly with the propagation of CSD. Soon after CSD propagation, the diameter recovered and thereafter it remained at the baseline level or decreased by 5% to 30% until the end of the measurement. In arteriole B, the diameter was limitedly changed at two

![Fig. 7 Maps of (a) regional hemoglobin concentration (rCHb) and (b) tissue oxygen saturation (StO₂) before and after LISW application for rat no. 3. The bottom left images in (a) and (b) are corresponding reflectance images at 520 nm before LISW application, showing identification of arterioles, venules, and parenchyma as “a,” “v,” and “p,” respectively. Pixel sizes in the images are 455 × 418.](https://www.spiedigitallibrary.org/journals/Journal-of-Biomedical-Optics)
positions (B-1 and B-2) over a period of ∼50 min after LISW application [Fig. 5(b)]. At the other position (B-3), however, the diameter decreased by ∼20% immediately after LISW application but transiently recovered. During CSD propagation, the diameter decreased again by ∼30%, after which the diameter gradually recovered to the baseline. In pial venules C and D, on the other hand, the diameters decreased by ∼45% and ∼25%, respectively, over a period of ∼1 min soon after LISW application and then both increased almost concomitantly with the propagation of CSD over a period of ∼1 min. Immediately after the CSD propagation, they decreased again by 20% to 30% from the baseline over a period of ∼1 min and then started to recover [Figs. 5(c) and 5(d)]. Interestingly, the diameters started to increase from the baseline at ∼7 min after LISW application and reached 120% to 155% at ∼20 min; the smaller venules showed larger variations in their diameters. Thereafter, the diameters gradually decreased in smaller venule C but showed little changes in larger venule D.

On the basis of the time courses of vessel diameters shown in Fig. 5, six characteristic time points were determined: $t_0$, before LISW application; $t_1$, immediately after LISW application; $t_2$, during CSD propagation; $t_3$, immediately after CSD propagation; $t_4$, ∼20 min post-CSD; and $t_5$, ∼50 min post-CSD [see arrowheads at the top of Fig. 5(a)]. Figure 6 shows changes in diameters of pial arterioles and venules relative to the baseline diameters (values before LISW application, $t_0$) at these time points for all four rats examined. For this analysis, we randomly selected a total of 14 positions from 5 arterioles for type A (running anterograde to the CSD propagation direction), 8 positions from 3 arterioles for type B (running perpendicularly to or retrograde to the CSD propagation direction), 17 positions from 8 venules for type C (with a diameter less than 80 μm), and 10 positions from 7 venules for type D (with a diameter of 80 μm or more); all of the observable arterioles and venules were evaluated. The mean diameters before LISW application were 48 ± 4 μm ($n = 14$, mean ± SEM) and 44 ± 5 μm ($n = 8$) for type A and type B arterioles and 53 ± 4 μm ($n = 17$) and 150 ± 13 μm ($n = 12$) for type C and type D venules, respectively. Unique responses depending on the vascular type were seen in all rats examined. The diameter of type A arterioles showed a significant increase at $t_2$ and the percent increase was 59 ± 9% (mean ± SEM). For type B arterioles, however, the diameter did not show a significant change at $t_2$ but was significantly decreased at $t_3$ (∼10 ± 3%). For both C and D types of venules, the diameters were significantly decreased at $t_1$ (−22 ± 4% and −20 ± 2%, respectively). After the propagation of CSD, both C and D types of venules showed significant increases at $t_4$ (34 ± 4% and 25 ± 3%, respectively). Vascular responses to the second consecutive CSD, which were observed only for rat no. 4, were excluded from analysis in this study.

### 3.3 Regional Hemoglobin Concentration and Tissue Oxygen Saturation

Figure 7 shows false-color maps of the distributions of $rCH_b$ and $StO_2$ before and after LISW application (rat no. 3; see also Fig. 8 for $StO_2$). The maps clearly show site-dependent variations in $rCH_b$ and $StO_2$ among the arterioles, venules, and parenchyma. Immediately after LISW application ($t = 1$ min 11 s), both $rCH_b$ and $StO_2$ decreased in almost all regions; their decreases in the parenchyma were remarkable and pial venules showed distinct constriction. Thereafter, $rCH_b$ and $StO_2$ started to increase in the right side region ($t = 1$ min 21 s), and the regions with increased $rCH_b$ and $StO_2$ were moving to the left side, corresponding to the propagation of CSD (until $t = 2$ min 29 s), indicating hyperemic responses to CSD (see Fig. 3). After the CSD propagation, $rCH_b$ and $StO_2$ started to drastically decrease ($t = 3$ min 26 s). Thereafter, $rCH_b$ tended to gradually recover but still continued to be lower than the values before LISW application in the parenchymal region ($t = 66$ min 37 s). On the other hand, $StO_2$ remained very low in all regions; the values in the parenchyma and venules were especially low until ∼40 min post-CSD and those in the parenchyma at ∼1 h post-CSD were still lower than the values before LISW application.

Figure 9 shows changes in $rCH_b$ and $StO_2$ in the parenchyma, arterioles, and venules at the six characteristic time points shown in Figs. 5 and 6 for all four rats examined, for which a total of 68 ROIs, including 24 ROIs in the parenchyma, 20 ROIs in arterioles running anterograde to the CSD propagation direction, and 24 ROIs in venules, were sampled. Immediately after LISW application (at $t_1$), $rCH_b$ was significantly decreased in all of the regions, by ∼8% in the parenchyma, by ∼6% in arterioles, and by ∼8% in venules from the baseline levels. During CSD propagation (at $t_2$, light blue shadow), $rCH_b$ was significantly increased in the parenchyma by ∼21% and in arterioles by ∼22%, and then immediately after CSD propagation (at $t_3$), $rCH_b$ was significantly decreased in the parenchyma, arterioles, and venules by ∼6%, by ∼4%, and by ∼8%, respectively. At $t_4$ and $t_5$ (more than ∼20 min post-CSD), while $rCH_b$ in the arterioles at $t_5$ was significantly decreased, $rCH_b$ had recovered to the baseline levels in the parenchyma and venules and did not change significantly. Corresponding to these $rCH_b$ changes, $StO_2$ also changed, but its behavior was more complex. At $t_1$, $StO_2$ was significantly decreased in all of the regions by ∼29% in the parenchyma, by ∼4% in arterioles, and by ∼14% in venules from the baseline levels, but at $t_2$, it was significantly increased in these regions by ∼64%, ∼14%, and ∼19% from the baseline levels, respectively. At $t_3$, $StO_2$ was significantly decreased by ∼46% in the parenchyma, by ∼10% in arterioles, and by ∼16% in venules from the baseline levels. After the CSD propagation (at $t_4$ and $t_5$), $rCH_b$ recovered.
Fig. 9 Regional hemoglobin concentration (rCHb), tissue oxygen saturation (StO$_2$), and arterial oxygen saturation measured at the hindlimb before and after LISW application for four rats examined; description for each time point is shown at the bottom of the graphs. (a–c) rCHb in the parenchyma (a), pial arterioles (b), and pial venules (c). (d–f) Corresponding StO$_2$ in the parenchyma (d), pial arterioles (e), and pial venules (f). (g) Arterial oxygen saturation (SpO$_2$) simultaneously measured at the hindlimb. The horizontal dashed line in each graph shows the baseline before LISW application. Numbers of RoIs used for analysis in rat no. 1, no. 2, no. 3, and no. 4 are as follows: 6, 6, 6, and 6 for parenchyma; 4, 6, 4, and 6 for arterioles; 6, 6, 6, and 6 for venules, respectively. rCHb levels in the parenchyma (a) are presented as medians with ranges; Dunn’s multiple comparisons test was performed. Other data are presented as means ± SD; Dunnett’s multiple comparisons test was performed. Asterisks (*/p < 0.05, ***/p < 0.01, *****/p < 0.0001) indicate a statistically significant difference.
to the baseline levels in almost all of the regions as described above, but SpO\text sub\text 2 still remained at significantly low levels, indicating a unique abnormal hemodynamic condition. Figure 9(g) shows arterial oxygen saturation (SpO\text sub\text 2) measured at the hindlimb simultaneously with rCH\text sub\text 2 and StO\text sub\text 2 for the cortex. It should be noted that even after LISW application, SpO\text sub\text 2 showed no significant changes at any time point, suggesting that the above-mentioned drastic changes in rCH\text sub\text 2 and StO\text sub\text 2 were local phenomena in the brain.

On the basis of the data of the arteriolar and venular SpO\text sub\text 2 described above, we evaluated OEF by the equation:

\[
\text{OEF} = \frac{1 - \text{StO}_2\text{,v} / \text{StO}_2\text{,a}}{\text{StO}_2\text{,v} / \text{StO}_2\text{,a}},
\]

where StO\text sub\text 2\text sub\text a and StO\text sub\text 2\text sub\text v are arteriolar and venular SpO\text sub\text 2, respectively. Mean values of StO\text sub\text 2\text sub\text a and StO\text sub\text 2\text sub\text v at each time point were obtained for all the RoIs described above, and OEF was then calculated for each rat. Figure 10 shows OEF values at all the six time points. The mean baseline OEF value (at t\text sub\text 0) was 0.26, and immediately after LISW application (at t\text sub\text 1), the value increased to 0.34. During CSD propagation (at t\text sub\text 2, hyperemic response to CSD), the mean OEF decreased to 0.22, but in the post-CSD period (from t\text sub\text 3 to t\text sub\text 5), the value increased again and reached to the maximum of 0.40 (at t\text sub\text 5, ~20 min post-CSD).

4 Discussion

Real-time observation of cerebral responses to a shock wave is crucial for understanding the pathophysiology and mechanisms of bTBI, because the initial events caused by a shock wave can trigger cascading processes leading to pathological consequences in the chronic phase. In the present study, we revealed the vascular type-dependent hemodynamic changes in the cortex from the moment of LISW application to the rat brain based on a multispectral imaging approach. In this model, we observed unique time-dependent changes in vascular diameters, rCH\text sub\text 2, StO\text sub\text 2, and OEF. It would be useful to analyze observed changes in the following four time regimes.

Phase I (t = t\text sub\text 1): immediately after LISW application (before the start of CSD, for ~1 min) in which there is immediate oligemia/hypoxemia with distinct venular constrictions.

Phase II (t = t\text sub\text 2): during CSD propagation (for ~2 min) in which there is a transient distinctive hyperemic response to CSD, i.e., increases in parenchymal rCH\text sub\text 2 and StO\text sub\text 2 with arteriolar dilatation.

Phase III (t = t\text sub\text 3): short-term post-CSD regime (for <20 min) in which there is oligemia/hypoxemia with arteriolar constriction.

Phase IV (t = t\text sub\text 4 ~ t\text sub\text 5): long-term post-CSD regime (for >40 min) in which there is long-lasting hypoxemia with neither oligemia nor arteriolar constriction but with remarkable increases in OEF (oxygen supply-demand mismatch).

The responses in phase I, i.e., direct, immediate cortical vascular and hemodynamic responses to a shock wave, are characterized by constrictions of both capillaries and venules; the parenchymal decrease in rCH\text sub\text 2 indicates capillary constrictions. It has been reported that diameters of cerebral capillaries are controlled by pericytes, and that capillary constriction is caused by contraction of pericytes under neuropathological conditions. It has also been shown that the active constriction of veins in acute brain injury is mostly due to contraction of pericytes. Thus, the immediate oligemia with venular constrictions observed in phase I is attributable to contractions of pericytes or otherwise to other contractile mural cells covering the capillaries and venules. It has been shown in a mouse ischemia-reperfusion model that generation of oxygen and nitrogen radicals, e.g., peroxynitrite, caused contraction of pericytes on microvessels, resulting in capillary constriction. Although the type of insult is different from ours, a similar mechanism might be involved. This is because overpressure of a shock wave can increase intracellular calcium (Ca\textsuperscript{2+}), resulting in activation of endothelial nitric oxide synthase (NOS) and hence an increase in nitric oxide (NO). Meanwhile, an intracellular increase in Ca\textsuperscript{2+} can stimulate the generation of reactive oxygen species (ROS) in mitochondria. The main ROS produced in mitochondria is known to be superoxide anion (O\textsuperscript{2−}), and O\textsuperscript{2−} and NO can efficiently react, producing peroxynitrite. We speculate that peroxynitrite is produced in the present model, leading to contraction of pericytes. In phase I, we also observed increased light scattering signals in the large venules during their constriction (Fig. 3). The venular constriction would involve cellular contractions in the vessel wall, which can result in an increase in local concentrations of macromolecules such as membrane lipid and collagenous substances, which are light scatterers. Thus, the scattering increase in the constricted venules would be associated with structural changes of those light scatterers.

Before starting phase II (CSD propagation), rCH\text sub\text 2 and StO\text sub\text 2 in the cortex had already been reduced with pial venular constriction due to the direct effects of the shock wave, as described above. Under normoxic conditions, massive hyperemia occurs in response to CSD, compensating the energy consumption used for restoration of imbalanced ions caused by CSD. In pathological conditions, such as mildly or moderately ischemic penumbra around a focal infarction, on the other hand, impaired cerebral blood flow (CBF) and oxygen metabolism can transform the CBF response from hyperemia to a hypoperfusion-dominant reaction, exacerbating tissue injury.
wave-induced CSD we observed showed the opposite response; distinct hyperemia occurred even under the initial hypoperfusion condition. This is probably due to the massive production of NO, which is a well-known vasodilator, from endothelial cells by the shear force caused by the shock wave. It should also be noted that a hyperemic response to CSD was evidently observed in arterioles running in the anterograde direction to the CSD propagation direction (Fig. 4). This would be associated with the phenomenon called vascular conduction of dilatation. It has been demonstrated that certain vascular stimuli, e.g., local application of KCl, induced rapid longitudinal conduction of vasomotor signals, resulting in propagation of the dilatory response. As one possible mechanism for the vascular conduction, electronic spread of membrane potential changes through gap junctions between vascular cells has been suggested.

The post-CSD oligemia/hypoxemia with arteriolar constriction observed in phase III is similar to the observation in our previous study based on fiber-based DRS. In that study, post-CSD hypoxemia was observed with oligemia during the whole period of measurements (≥1 h). In the present study, however, vascular constriction and oligemia recovered within ~20 min post-CSD in phase III, but hypoxemia continued into phase IV. This difference is attributable to the different observation depths by the different methodologies and wavelengths used in the two studies; fiber-based DRS in the visible to near-infrared region (previous study) and CCD-based multispectral imaging in the visible region (present study). Observable tissue regions can be estimated as being the full thickness of the cortex (layers I-VI) in the previous study and the superficial cortical layer, i.e., mainly layer I, in the present study. In the rat cerebral cortex, vascular anatomy changes dramatically depending on cortical depth, e.g., capillary density increases as a function of depth, while larger-diameter arterioles and venules are dominantly located on the cortical surface. It was shown by depth-resolved laser Doppler flowmetry for rat brains that after a hyperemic response to CSD, the mean flow values for pial arterioles, venules, and other small vessels located immediately below or at the surface of the brain soon recovered to the baseline values, while the flow values at a depth of ~500 μm in the cortex were decreased for a while (approximately >8 min). Although the exact mechanisms are not clear, we assume that post-CSD persistent hypoxemia observed in the deeper cortical region results from oligemia associated with deeply located arterioles/venules and capillaries, while that observed in the superficial region is caused by the mechanisms discussed below. The hypoxemia that persisted even after recovery of blood volume with distinct venular dilatation in phase IV indicates a hypoxic state and/or a hypermetabolic state, which seems to be unique to shock wave-related CSD responses. This can be caused by inadequate oxygen supply from blood vessels and/or excessive oxygen consumption. The mechanism for the decreased StO2 even in arterioles [Fig. 9(e), at t4 and t5] is not clear, but it may be associated with an acidic tissue condition after CSD propagation. It has been shown that extracellular lactate was increased and pH was decreased for more than 30 min after KCl-induced CSD onset in cats. Such an acidic condition can also occur in the present model after CSD propagation. In addition, the hypoxic state with restored rChb in this phase suggests that the tissue was supplied with glucose under the condition of insufficient oxygen delivery, accelerating lactate accumulation and hypoxic brain acidosis. The lowered pH decreases hemoglobin affinity for oxygen as explained by the hemoglobin-oxygen dissociation curve, resulting in decreased arteriole StO2.

Regardless of the decreased StO2 in phase IV, cellular oxygen metabolism in cortical tissue may increase in this phase, which is supported by the remarkable increases in OEF at t4 and t5 (Fig. 10). The initial oligemia/hypoxemia directly caused by the shock wave itself (Phase I) and the subsequent CSD (Phase II) and post-CSD hypoxemia/oligemia (Phase III) can all be regarded as a process that imposes metabolic impacts on cortical tissue. During these periods, oxidative and nitrosative stresses might have occurred successively, because CSD can be accompanied by NO production from neuronal NOS and increase in intracellular Ca2+, leading to the generation of ROS. These reactions and products can increase cellular metabolism, e.g., specific protein synthesis, oxidative phosphorylation and ATP formation, for recovery of cellular functions, resulting in excessive oxygen consumption. Actually, OEF increased significantly in post-CSD Phase IV (at t4 and t5); the mean baseline OEF value (0.26) increased to 0.40 at ~20 min post-CSD (at t4). It is interesting to know that a similar OEF increase has also been reported in ischemic penumbra in focal cerebral ischemia model. To cope with such an increased energy demand, the cortical tissue would have to struggle with energy production because capillary density is relatively low in the cortical surface region. From a biomechanical point of view, the cortical surface region just beneath the skull might encounter a specific type of injury resulting from cavitation because a shock wave is partly reflected at the interface between the cerebral spinal fluid and the skull.

To validate the hypotheses for the mechanisms of hemodynamic responses to a shock wave described above, we plan to perform in vivo imaging of calcium ions and NO, which are likely to be key players in the superacute events. Assessments of tissue pH and pressure characteristics just beneath the skull are also important parts of our future work. We are also currently concerned about histopathological consequences of the unique vascular and hemodynamic responses to a shock wave and the concomitant CSD observed in this study. Analyses of the related possible secondary brain damage, such as oxidative stress, blood brain barrier disruption, apoptosis, and neuroinflammation, are ongoing at present.

In summary, we here reported what we believe to be the first real-time observation of vascular type-dependent responses of hemodynamics and vessel morphology to a shock wave in the rat brain by using multispectral reflectance imaging. It was shown that hemodynamic abnormalities were caused solely by a shock wave and they were exacerbated by the occurrence of CSD. Vasoconstriction was more evident in venules than in arterioles, suggesting pericyte alteration due to a shock wave. After CSD propagation, a persistent hypoxic state and/or hypermetabolic state (an oxygen demand-supply mismatch) was observed. These findings may give new insights into the understanding of the blast TBI mechanisms.

5 Appendix

The objective of this appendix is to compare the results of calculation of hemodynamic changes obtained by the multiple regression analysis described in the main text and those obtained by the conventional modified Lambert–Beer analysis. In comparison, we used the same image datasets obtained at
Fig. 11 Changes in concentrations of (a) oxygenated, (b) deoxygenated, and (c) total hemoglobins in the rat brain exposed to an LISW (rat no. 3) obtained by multiple regression analysis aided by Monte Carlo simulation. Each time point is the same as Fig. 7.
Fig. 12. Changes in concentrations of (a) oxygenated, (b) deoxygenated, and (c) total hemoglobins in the rat brain exposed to an LISW (rat no. 3) obtained by the conventional modified Lambert–Beer law. Each time point is the same as Fig. 7.
The modified Lambert–Beer law relates a measured reflectance intensity change to changes of chromophore concentrations. Assuming that oxygenated and deoxygenated hemoglobins are the only chromophores in the tissue of interest, the relationship between the measured reflectance intensity change and changes in $C_{HbO}$ and $C_{HbR}$ is described as the following equation:

$$\Delta A(\lambda) = \{\epsilon_{HbO}(\lambda) \Delta C_{HbO} + \epsilon_{HbR}(\lambda) \Delta C_{HbR}\} D_o(\lambda), \tag{10}$$

where $\Delta A(\lambda) = \log(R_0/R_t)$ is the attenuation at each wavelength, $R_0$ and $R_t$ are the measured reflectance intensities at time 0 (baseline) and at time $t$, $\Delta C_{HbO}$ and $\Delta C_{HbR}$ are concentration changes of HbO and HbR, respectively, and $\epsilon_{HbO}$ and $\epsilon_{HbR}$ are the respective molar extinction coefficients. Equation (10) can be solved for $\Delta C_{HbO}$ and $\Delta C_{HbR}$ using a least-square approach. The differential pathlength factor, $D_o(\lambda)$, accounts for the fact that a photon at each wavelength travels a slightly different optical path through the tissue due to the wavelength-dependent scattering and absorption coefficients in the tissue, and it can be estimated using the approach developed by Kohl et al.\textsuperscript{22} on the basis of MCS of light propagation in a homogeneous tissue model. The wavelength-dependent reflectance, $R_{\mu_0}(\lambda)$, can also be calculated using MCS to determine $D_o(\lambda)$.\textsuperscript{22} In the analysis here, the baseline reduced scattering coefficient, $\mu_0$ was assumed to be 12 to 16 cm$^{-1}$ at 520, 540, 560, 580, and 600 nm\textsuperscript{26} and was assumed to be constant with time according to the literature.\textsuperscript{22,29}

Figures 11 and 12 show images of the changes in $C_{HbO}$, $C_{HbR}$, and $C_{HbT}$ obtained by the multiple regression analysis and the conventional modified Lambert–Beer analysis, respectively (rat no. 3). Time courses of changes in $C_{HbO}$, $C_{HbR}$, and $C_{HbT}$ shown in Fig. 11 are consistent with the time-dependent changes in regional hemoglobin concentration ($rC_{Hb}$) and tissue oxygen saturation ($SO_2$) shown in Fig. 7. It should be noted, however, that some regions at the vessel edges show extremely high or low values, which is attributable to changes in these vessels’ diameters. Figure 12 shows the corresponding results obtained by the conventional modified Lambert–Beer analysis. The tendencies similar to those in Fig. 11 can be seen in Fig. 12. However, the values of changes in Fig. 12 are slightly smaller than those in Fig. 11. This may be due to the different assumption regarding the scattering coefficient; in the modified Lambert–Beer law analysis, the scattering coefficient was treated as constant. The authors plan to conduct more comprehensive comparison between the two methods.

Disclosure

The authors have no conflicts of interests related to this work.

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References


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