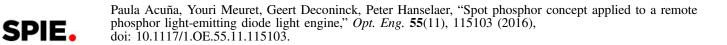
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Paula Acuña Youri Meuret Geert Deconinck Peter Hanselaer



### Spot phosphor concept applied to a remote phosphor light-emitting diode light engine

Paula Acuña,<sup>a,\*</sup> Youri Meuret,<sup>a</sup> Geert Deconinck,<sup>b</sup> and Peter Hanselaer<sup>a</sup>

<sup>a</sup>KU Leuven ESAT/Light and Lighting Laboratory, Gebroeders de Smetstraat 1, Ghent 9000, Belgium <sup>b</sup>KU Leuven ESAT/ELECTA, Kasteelpark Arenberg 10 bus 2445, Leuven 3001, Belgium

**Abstract.** Although remote phosphor technology outperforms conformal phosphor technology for midpower applications, one of the limiting factors due to its impact on the total cost is the amount of phosphor required. Furthermore, an important loss mechanism in remote phosphor light-emitting diode (LED) technology is the reabsorption of recycled, downconverted light by the phosphor. An obvious solution to this issue is to enable a light path for the converted light, such that further interactions with the phosphor element are avoided. We propose a spot phosphor concept to achieve this goal. To explore this configuration, a simulation model of a phosphor element is devised and validated. The optical input parameters are based on experimental data and the application of the inverse adding-doubling method. The resulting configuration, along with a long-pass filter, is shown to be a potential solution for reduction of phosphor usage. The moderate decrease in the light extraction ratio (LER) when applying the spot concept is partly attributed to the losses in the secondary optics needed to narrow the LED beam; the combination of the spot concept configuration with a directional light source such as a laser diode is shown to be a powerful combination for the enhancement of the LER. © *The Authors. Published by SPIE under a Creative Commons Attribution 3.0 Unported License. Distribution or reproduction of this work in whole or in part requires full attribution of the original publication, including its DOI.* [DOI: 10.1117/1.OE.55.11.115103]

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#### 1 Introduction

It has been demonstrated that the light extraction efficiency of a remote phosphor light-emitting diode (LED) light engine in a cylindrical configuration can only be moderately increased by varying the optical and geometrical parameters.<sup>1,2</sup> After examining the interaction of yellow light with a phosphor plate, both Acuña et al.<sup>3</sup> and Liu et al.<sup>4</sup> identified that nearly 10% of phosphor-converted light recycled back toward the phosphor is absorbed by this phosphor element. Furthermore, rare earth-based phosphor is one of the components that drive up the cost of solid-state based lighting. According to the manufacturing roadmap for solid-state lighting of 2014,<sup>5</sup> phosphors represent 8% and 18% of the total manufacturing costs for high-power and medium-power LEDs, respectively. In addition, the contribution of the phosphor to the total cost will become higher, since envisaged progress in solid-state technology addresses mainly the cost reduction of different components, but not that of the phosphor. To address the loss mechanisms and to reduce the usage of phosphor, the most intuitive solution is to avoid any reinteraction of the emitted downconverted light with the phosphor element.

Attempts to address the light extraction efficiency and color uniformity of white phosphor-converted LEDs by changing the shape of the phosphor have been restricted to single-chip packages.<sup>6–10</sup> However, the architecture of a remote phosphor LED module differs considerably regarding the size, geometry, and construction of the light source configuration, which permits completely different trends and

behavior. In this paper, a remote phosphor concept in which the phosphor surface is only a portion of the total emitting surface of the remote phosphor module is proposed (see Fig. 1). The part of the emitting surface not covered by phosphor remains transparent. As a consequence, the pump source angular emission must be reduced as well. In this way, the light emitted by the pump source hits exclusively the module's exit aperture, which is coated with phosphor. Both the blue and yellow backward-scattered light emitted by the phosphor reach the mixing cavity surrounding the pump source and are redirected toward the module's full exit aperture, which is only partly covered with phosphor. Hence, the reabsorption of yellow light in the phosphor can be suppressed. This configuration will be referred to throughout the text as the spot concept.

To find the physical dimensions and phosphor concentration that produced white light with the maximum light extraction ratio (LER) for a certain color point, a simulation model was developed and validated, as presented in Sec. 2. Using the validated model, the spot phosphor concept activated by blue-emitting LEDs is explored in Sec. 3. As an alternative to the use of secondary optics to collimate the emission of the blue-emitting LEDs, the performance of the spot phosphor concept with blue-emitting laser diodes (LDs) is presented in Sec. 4. Results and conclusions are given in Sec. 5.

#### 2 Simulation Model

The optical characteristics of a phosphor element comprise its luminescence properties, surface scattering, and volume scattering. The luminescence properties are its excitation and emission spectra and quantum efficiency. The surface scattering is characterized by the bidirectional scattering

<sup>\*</sup>Address all correspondence to: Paula Acuña, E-mail: paula.acuna@kuleuven .be

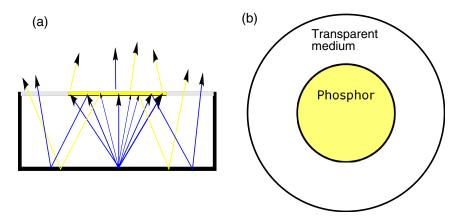


Fig. 1 Spot phosphor concept: (a) front and (b) top view.

distribution function (BSDF), and the volume scattering is defined by the bulk scattering parameters. The BSDF is defined as the ratio of the differential spectral radiance of the sample at a particular viewing angle to the differential spectral irradiance on the sample from a particular incident angle

$$q_e = \frac{dL_{e,s}(\theta_i, \phi_i, \theta_s, \phi_s)}{dE_{e,i}(\theta_i, \phi_i)}.$$
(1)

The bulk scattering parameters comprise the absorption coefficient, the scattering coefficient, and the phase function. (Bidirectional reflectance distribution function [BRDF] and bidirectional transmission distribution function are similarly defined for reflection and transmission.)

To characterize the optical characteristics of the phosphor used in the final concept, a phosphor element was considered that consisted of a YAG:Ce-silicone oil solution with a phosphor concentration of 50 mg/ml immersed in a glass cuvette with a gap between the glasses of 1.1 mm.<sup>11,12</sup> The refractive index of the silicone is determined by the phenyl-methyl concentration. Given that the phenyl-methyl concentration is the same for both the silicone oil and the silicone encapsulant, no significant change in the optical behavior is expected.<sup>13</sup> Due to the smooth surface of the glass, the surface scattering is avoided and the radiant intensity of the slab under irradiation can be interpreted as a combination of the Fresnel reflectance and transmittance in the specular directions, the volume scattering and the fluorescence of the YAG:Ce particles diluted in the silicone oil.

The volume scattering optical parameters are obtained by solving the inverse problem of the radiative transfer equation (RTE) by means of the inverse adding-doubling (IAD) method applied to the experimental data of the transmission and reflection radiant intensities.<sup>14</sup> Two wavelengths have been chosen, 450 and 580 nm, corresponding to the excitation and emission spectrum regions. The reflection and transmission radiant intensity values are measured using a custom-made BSDF setup (see Ref. 15 for further details). The results for short- (450 nm) and long-wavelength (580 nm) illuminations are depicted in Figs. 2 and 3. From these data, the calculated values for the total reflectance *R*, total transmittance *T*, and the regular transmittance *T*<sub>r</sub> for both wavelengths 450 and 580 nm are calculated and indicated in the column "measurement" in Table 1.

The inverse problem solution to the RTE using the IAD has been performed using a MATLAB<sup>®</sup> code developed and validated by Leyre et al.<sup>11</sup> and updated by Córreia et al.<sup>16</sup> Using genetic algorithms, the best-fit between experimental and calculated results is found considering 256 quadrature points in the discretization. The Gegenbauer–Kernel was adopted as the phase function<sup>17</sup> that best fits to the experimental data

$$p_{GK}(\theta) = K \cdot [1 + g^2 - 2 \ g \cos(\theta)]^{-(\alpha+1)},$$
  

$$K = \alpha g \frac{(1 - g^2)^{2\alpha}}{\pi [(1 + g)^{2\alpha} - (1 - g^{2\alpha})]},$$
(2)

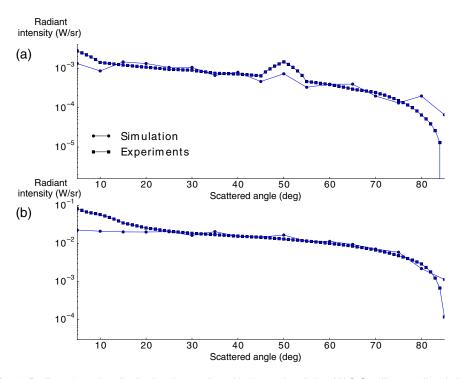
with g as the anisotropy factor,  $\theta$  as the angle to the normal, and  $\alpha$  as a fitting constant.

The IAD solution optimizing the match between the measurements and the simulation model based on a merit function considering both the shape and the absolute difference values of the angular-resolved transmission and reflection. The shape and the absolute differences are quantified by means of the normalized-cross-correlation function (NCC) [Eq. (3)] and the normalized-root-mean-squareerror (NRMSE) [Eq. (4)], respectively, where  $m, c, \theta_i, I,$ N and  $\overline{\cdot}$  stand for measured, calculated, angle to the normal, radiant intensity, number of angles measured, and average, respectively

$$\operatorname{NCC} = \frac{\sum_{\theta_i} \{ I_m(\theta_i) - \overline{I_m(\theta_i)} \cdot [I_c(\theta_i) - \overline{I_c(\theta_i)}] \}}{\left\{ \sum_{\theta_i} [I_m(\theta_i) - \overline{I_m(\theta_i)}]^2 \sum_{\theta_i} [I_c(\theta_i) - \overline{I_c(\theta_i)}]^2 \right\}^{1/2}},$$
(3)

$$NRMSE = \frac{\sum_{\theta_i} \sqrt{[I_c(\theta_i) - I_m(\theta_i)]^2}}{\max(I_c, I_m)N}.$$
(4)

The NCC varies between 0 and 1, 1 being a perfect match and 0 absence of correlation between simulation and measurement. The NRMSE gives values higher or equal to zero, zero being a perfect match. Finally, the relative error of the integrated quantities, i.e., total reflectance and total transmission, is also considered by the merit function, which is defined as



**Fig. 2** Radiant intensity distribution in semilogarithmic scale of the YAG:Ce-silicone oil solution 50 mg/ml illuminated with a collimated beam with central wavelength 450 nm, FWHM 10 nm. Circleand square-lines represent the measurement and the simulation, respectively. Transmission and reflection distribution are illustrated in (a) and (b), respectively.

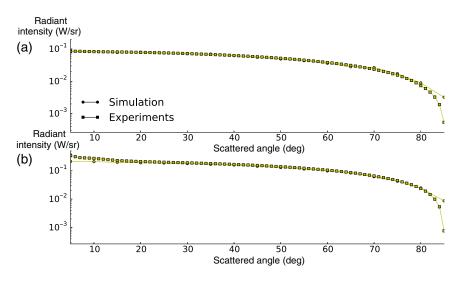


Fig. 3 Radiant intensity distribution in semilogarithmic scale of the YAG:Ce-silicone oil solution 50 mg/ml illuminated with a collimated beam with central wavelength 580 nm, FWHM 10. Circleand square-lines represent the measurement and the simulation, respectively. Transmission and reflection distribution are illustrated in (a) and (b), respectively.

Table 1	Total r	eflectance	and	transmittance	values	for th	ne samp	le und	ler test	
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	At 450	) nm	At 580	) nm
Parameter	Measurement	Simulation	Measurement	Simulation
Regular transmission	$2.08 \times 10^{-4}$	$1.99 \times 10^{-4}$	$7.04 \times 10^{-4}$	$6.57 \times 10^{-4}$
Total transmission	0.1804	0.1828	0.6014	0.5790
Total reflectance	0.0752	0.1164	0.3964	0.3965

 
 Table 2
 Bulk scattering and absorption parameters obtained from the IAD calculation for the phosphor-silicone oil 50 mg/ml solution with physical thickness 1.1 mm.

	At 450 nm	At 580 nm
Scattering coefficient $\mu_s$ (mm <sup>-1</sup> )	7.74	6.56
Absorption coefficient $\mu_a$ (mm <sup>-1</sup> )	0.48	0.0026
Anisotropy factor $g$ (p.u)	0.97	0.95
α (a.u.)	0.15	0.2

$$MF = (1 - NCC_{tra}) + NRMSE_{tra} + \frac{T_c - T_m}{T_m} + (1 - NCC_{ref}) + NRMSE_{ref} + \frac{R_c - R_m}{R_m},$$
(5)

where the subindices tra and ref stand for transmission and reflection, respectively.

Angular-resolved results of reflection and transmission from the simulation are depicted for illumination with light at 450 and 580 nm in Figs. 2 and 3, respectively. The NCC and the RMSE quantifying the differences between experimental and simulation results are above 0.97 and below 5.4%, respectively. The integrated values R, T, and  $T_r$  from the simulation are tabulated in Table 1 for comparison with the measured values.

The obtained parameters are summarized in Table 2. These are comparable with those reported for the same phosphor material in Refs. 4, 14, and 18. These absolute values have been applied as scale factors for the relative spectral curves.<sup>19</sup>

The quantum yield, defined as the ratio of converted photons to the absorbed photons, is obtained from the manufacturer Phosphor Technologies and is shown in Fig. 4. Both the absolute quantum yield value at 450 nm and the emission spectrum are measured with a custom-made integrating sphere (see Ref. 20 for further details). The resulting quantum yield value at 450 nm equates to 0.88. A simulation model is implemented along with the pump source and the mixing cavity to evaluate the performance of the spot phosphor concept.

Both the scattering and the luminescent properties of the phosphor are modeled and simulated wavelength per wavelength using the commercial ray-tracing software LightTools.

#### 3 Spot Phosphor Concept Activated by Blue Light-Emitting Diodes

The remote phosphor configuration that we consider consists of a cylindrical mixing cavity of 50-mm radius and 20-mm height, with blue-emitting LEDs located at the base and the phosphor element with the same radius as the mixing cavity located at the top. The inner surface of the mixing cavity consists of Mylar (polyethylene terephthalate with a total reflectivity value of 98.5%, and a rather diffuse pattern (98% is diffuse, 2% specular). The inner surface is modeled using the BRDF. The pump source is a luxeon royal blue LED from Lumileds with central wavelength at 450 nm, with emitting surface of 1 mm × 1 mm and rather Lambertian emission. The optical emission of the pump source has been set to 1 W of radiant power and is modeled by using a ray file with five million rays.

In the baseline configuration, the exit aperture is totally covered with a uniform phosphor element, while the spot phosphor concept has a reduced phosphor surface. As a consequence, in the spot concept, the pump source angular emission must be reduced as well, while part of the emitting surface remains uncovered by phosphor (see Fig. 1).

However, LEDs typically emit a Lambertian intensity distribution. For this reason, the use of secondary optics to confine the flux to a smaller solid angle is required. According to the *étendue* conservation principle,<sup>21</sup> the product of the spatial extent and the angular extent remains constant in an optical system.

The *étendue* of a naked LED is calculated as the product of the small emitting area and the large angular extent. If the angular extent has to be decreased to hit only the small phosphor area, the exit aperture of the secondary optical element should be larger than the LED emitting area.

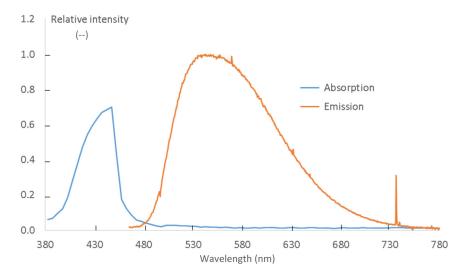


Fig. 4 Relative absorption and emission spectra of the phosphor powder YAG:Ce in a oil-solution of 50 mg/ml measured with the quantum efficiency setup.

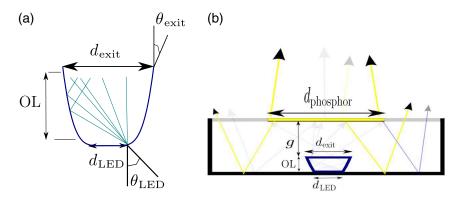


Fig. 5 (a) CPC geometry and (b) its application in the spot architecture.

The implementation of the secondary optics, in this paper, is based on reflective optics, namely an aluminum compound parabolic concentrator (CPC) (see Fig. 5), with a specular reflectance of 95%. Following the *étendue* conservation, the CPC exit diameter  $d_{\text{exit}}$  can be defined in terms of the CPC input diameter  $d_{\text{LED}}$  and the maximum deflection angle at the exit  $\theta_{\text{exit}}$  as

$$d_{\text{exit}} = \frac{d_{\text{LED}}}{\sin(\theta_{\text{exit}})}.$$
(6)

The CPC height can be expressed as<sup>21</sup>

$$OL = \frac{d_{LED}/2 \cdot [1 + \sin(\theta_{exit})] \cdot \cos(\theta_{exit})}{\sin^2(\theta_{exit})}.$$
(7)

When optimizing the geometry of the spot architecture only one constraint should be satisfied. The height of the CPC "OL" plus the gap between the CPC exit and the phosphor g must be equal to the height of the mixing cavity, i.e., 20 mm. In addition, the CPC exit surface projection at a distance g should equate to the phosphor surface, hence, the deflection angle  $\theta_{exit}$  must comply with the following equation:

$$d_{\text{phosphor}} = d_{\text{exit}} + 2 \cdot \tan(\theta_{\text{exit}}) \cdot g.$$
(8)

The diameter of the phosphor coating has been reduced to 50 mm, the LED chip size ( $d_{\text{LED}}$ ) was set to 2.6 mm according to the silicone encapsulant on the emitting surface of the LED and the angular beam to 90 deg. The resulting dimensions of the CPC are specified in Table 3.

The merit functions used to evaluate a configuration are the LER, the luminous efficacy of radiation (LEM), the luminous efficacy (LE), the correlated color temperature (CCT), the color rendering index (CRI), the color over angle (CoA) variations (Du'v'), and the relative phosphor mass. The CoA variation is quantified as the Du'v' at angle  $\theta$  with respect to  $\theta = 0$  deg and it is calculated according to the European Standard EN13032-4,<sup>22</sup> i.e., angular intervals of 22.5 deg and 2.5 deg in the vertical and horizontal plane, respectively. The max(Du'v') taken as reference is 0.004, according to the Energy Star program requirements.<sup>23</sup> The thickness of the phosphor element has been adjusted such that the CCT and the CRI are the same as for the baseline configuration, 6800 K and 75, respectively. However, for any reasonable thickness compared to the commercial archetypes, the

 Table 3
 Dimensions of the CPC in the remote phosphor module illuminated by a LED.

Dimension	Value
d <sub>phosphor</sub>	50 mm
d <sub>LED</sub>	2.6 mm
$\theta_{LED-max}$	90 deg
Maximum value $\theta_{\text{exit}}$	40
OL	4.0 mm
g	16.0 mm
d <sub>exit</sub>	4.1 mm

CCT was always higher due to the amount of blue light escaping through the noncoated surface after the recycling process. Therefore, the radius of the coated region was increased such that the CCT of 6800 K could be achieved with a more realistic thickness. The resulting radius and thickness were 38 and 6.0 mm, respectively. Considering a phosphor concentration of 50 mg/ml, the phosphor mass (*m*) results in 1066 mg. The reflectance of the LED emitting surface was set to 60%, according to the value reported by Cornelissen et al.<sup>24</sup>

A summary of the merit functions is presented in Table 4.

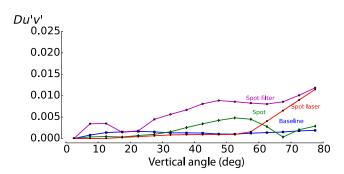
For the same color point, the spot concept LER is 3% lower than for the baseline, which is also reflected on the lower values of the LEM and LE. The phosphor mass is 32% larger than the baseline counterpart.

Since the phosphor element has become thicker than in the baseline, more yellow light is backward scattered after the first interaction compared to the baseline configuration. Also, as a consequence of the thicker phosphor element, the absorption losses due to limited quantum efficiency and Stokes shift after one interaction are larger in the spot concept (28%) than in the baseline (22%) configuration. The interaction of backward scattered blue light with the phosphor element is significantly less (35% of the baseline counterpart); however, the amount of recycled light is 20% larger for the spot concept.

The CoA with the spot concept is slightly larger than the baseline as a consequence of the nonuniformly coated exit surface (see Fig. 6). The Du'v' increases with the angle,

Parameter	Units	Baseline	Spot	Spot with filter	Spot laser
Electrical power	W	2.2 (2.8)	2.2	2.2	2.8
Radiant power (light source)	W	1 (1)	1	1	1
Radiant power (collimated beam)	W	0.95 (1)	0.95	0.95	1
LER	p.u	0.55 (0.70)	0.52	0.58	0.73
LEM	lm/W <sub>opt</sub>	187 (230)	181	187	237
LE	Im/W <sub>elec</sub>	84 (82)	81	85	83
ССТ	к	6808 (6783)	6774	6835	6828
CRI	_	75(69)	75	75	69
Thickness	mm	2.06 (1.15)	6.0	3.3	1.6
d <sub>phosphor</sub>	mm	100(100)	76	60	27
т	mg	808(143)	1066	629	42

 Table 4
 Merit figures and their values for the baseline and the spot configurations. The numbers between brackets are the baseline values obtained with the LD.



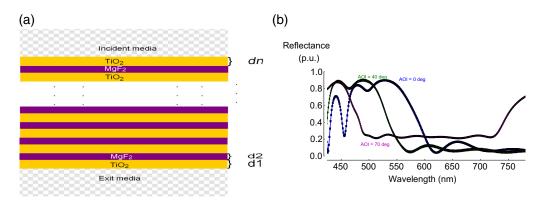
**Fig. 6** CoA variations in the u'v' space for the four tested configurations: baseline, spot LED activated, spot LED activated with filter, and spot-filter concept.

turning more bluish at larger angles as consequence of the blue light escaping through the noncoated region in the recycling process.

Although the spot concept seemed to be a logical approach to avoid reabsorption, it turns out that this approach does not provide a better performance and introduces various disadvantages. However, the results could be improved significantly by decreasing the phosphor thickness and using a long-wave pass thin-film filter on the phosphor noncoated region of the exit aperture.

The materials used for the thin-film filter were  $\text{TiO}_2$ and MgF<sub>2</sub>, with refractive indices of 2.346 and 1.384, respectively. The layers' stacking follows the formulation  $(0.5 HL0.5 H)^9$ , which is the typical stack of a long-wave pass filter. The merit functions for the layer thickness optimization were the cut-on wavelength (470 nm) and the incident angle (70 deg), according to the angular distribution of the rays incident onto the phosphor-free region of the exit aperture. For the design of the thin-film filter, the open source software Open Filters was used,<sup>25</sup> which uses the needle method as the optimization process. The scheme of the resulting thin-film filter and the spectral reflectance are illustrated in Figs. 7(a) and 7(b), respectively. The spectral reflectance is depicted for three angles of incidence (AOI), 0 deg, 40 deg, and 70 deg.

The results obtained with the spot concept combined with the filter are gathered in Table 4. The optical performance is



**Fig. 7** Long-wave pass filter based on thin-film coating of  $TiO_2$  and MgF<sub>2</sub>: (a) schematic representation and (b) spectral reflectance for AOI of 0 deg, 40 deg, and 70 deg.

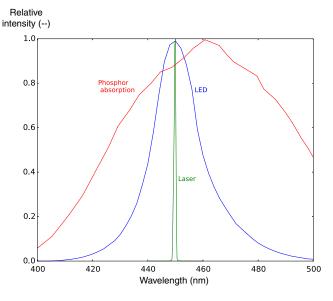


Fig. 8 Spectral distribution of a blue LED and a blue LD.

very comparable to the baseline's performance. The LER achieved with this configuration is 56%, very similar with the baseline. However, the amount of phosphor is significantly reduced to 77%. Regarding the Du'v' of the spot concept with filter in Fig. 6, it exhibits an increase with the angle, expecting to be more yellowish at larger angles as consequence of the filter.

#### 4 Spot Phosphor Concept Using a Blue Laser Diode

An alternative way to implement the spot concept without using additional optics to collimate the beam is to use an LD, characterized by a narrow angular beam.<sup>26</sup> A typical LD possesses an elliptical intensity distribution as a consequence of the active layer rectangular shape. The LD TB450B manufactured by Osram GmbH has been considered in our study. This LD typically emits light at 450 nm with an FWHM of 1.8 nm and has a maximum optical output of 1.6 W at a maximum angular divergence of 23 deg.<sup>27</sup> The LD TB450B emitting surface possesses a diameter of 2 mm.

In addition, the spectral power distribution is much narrower than that of LEDs, as can be seen in Fig. 8. Furthermore, the absence of efficiency droop constitutes another advantage, since the LD can be driven at a higher current without degrading the conversion efficiency.

In order to achieve a CCT of 6800 K and a CRI of 75, the phosphor thickness should, in this case, be only 1.55 mm compared to the 6 mm when using the LED-pumped spot concept. This is attributed to the narrow FWHM of the LD spectral power distribution (see Fig. 8). Less radiation away from the phosphor absorption peak (460 nm) implies less scattered light at short wavelengths, which increases the yellow-to-blue ratio and, thereby, decreases the CCT. The LD TB450B has been modeled based on the ray file provided by the manufacturer with a radiant power of 1 W. The optical and geometrical properties of the mixing cavity have been kept the same as in the previous configuration. For a distance between the LD and the phosphor layer of 20 mm and a maximum emitting angle of 35 deg, the minimum phosphor radius should be 14 mm.

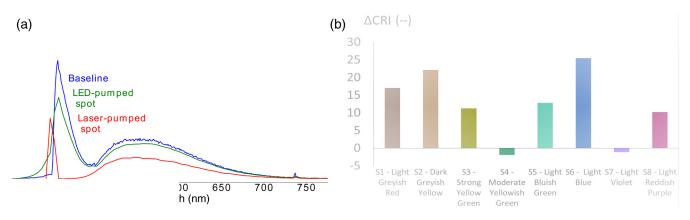
The simulation results of the laser-pumped spot configuration are summarized alongside those of the two previous configurations in Table 4. To make a fair comparison, the values within brackets in the column baseline represent the baseline concept but using the LD instead of the LED. The baseline pumped with LD has been tuned such that the same color point, i.e., 6800 K, is generated. The LER of the laser-pumped spot configuration is superior to both the baseline (LED based and LD based) and the LED-pumped spot concept values. This is attributed to the smaller phosphor-coated region at the exit plane, which reduces the opportunity for the recycled light to interact again with the phosphor element. Since the laser spectral emission is concentrated in a narrow band of wavelengths localized very close to the phosphor absorption spectrum peak, the absorption and subsequent emission at longer wavelengths are higher for the laser-pumped spot configuration, which result in the higher LEM. Currently, the luminous efficacy reported for a white phosphor-converted LD at a CCT of 6000 K is about 45 lm/W.<sup>28</sup> It is worth mentioning that the LD state of the art possesses a lower radiant efficiency than the LED for blue light, i.e., 35% compared to 45% for a driving current of 350 mA. Unfortunately, the lower radiant efficiency of the laser compared to the LED degenerates the good photoluminescent performance of the laser-pumped spot configuration to almost the same value as the baseline concept. However, this concept illustrates the potential for using LD when higher efficiencies will be reached in the future.

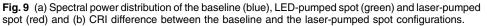
The advantage of the narrow spectral emission of the LD, however, constitutes a disadvantage from the CRI standpoint. The CRI of the laser-pumped spot configuration is 11 and 6 points below the baseline value and the LED-pumped spot configurations, respectively. This can be visualized via the spectral power distribution for all three configurations in Fig. 9(a). The narrow blue peak and the absence of power within the wavelength region 420 to 445 nm and 455 to 470 nm are not desirable but could be solved by adding other phosphors. Comparing the CRI for each test color sample, the largest difference appears for the color sample number 6, i.e., light blue, as depicted in Fig. 9(b).

The Du'v' variations for angles smaller than 60 deg are below the maximum defined in the energy star standard (Du'v' = 0.004),<sup>23</sup> as observed in Fig. 6 (left). This is originating from the concentration of blue light power toward the angle close to the normal, as expected from the laser angular emission.

An apparent disadvantage of concentrating the blue light onto a smaller phosphor area is that the temperature of the phosphor becomes much higher.<sup>29,30</sup> According to literature,<sup>31</sup> YAG:Ce-silicone white phosphor-converted LEDs can withstand phosphor luminance levels of up to 1e6 cd/m<sup>2</sup> without undergoing thermal quenching in the steady state. In our device, the luminance reaches a maximum value of 2.2e5 cd/m<sup>2</sup>, hence, no thermal quenching is expected to occur. For higher luminance levels, the thermal quenching present in the silicone-phosphors<sup>32</sup> is no longer an obstacle thanks to the development of single crystal phosphors<sup>33</sup> with higher thermal quenching temperatures.

Finally, it is worth mentioning that for the laser-pumped spot concept, the amount of phosphor usage drops to only





6% and 30% of the phosphor mass required for the baseline concept pumped with LED and LD, respectively.

#### **5** Conclusions

A simulation model of a phosphor element that considers both the surface and volume scattering, as well as the luminescence effects, have been implemented and validated with experimental data. Based on the simulation model, the performance of the spot concept architecture applied to a remote phosphor LED module has been evaluated and compared to the baseline. Results reveal that no significant improvement is achieved regarding the system's LER; however, there is a very significant reduction in phosphor usage when the spot concept is extended with a thin-film long-pass filter. The implementation of the spot concept with a more directional light source, such as a LD, shows high potential regarding luminous efficacy and phosphor reduction but needs to be investigated in more depth to improve CRI and angular color consistency.

#### Acknowledgments

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**Paula Acuña** received her BE and MSc degrees in electronics and electrical engineering, respectively, from the National University of Colombia, Bogota, in 2009 and 2011, respectively. Since 2011, she has joined the Light and Lighting Laboratory, Ghent, to pursue her doctoral studies in association with the Katholic University of Leuven. Her current research interests include solid-state lighting and lighting design with ray-tracing and photoluminescent materials.

**Youri Meuret** graduated in physics from Ghent University, Belgium, in 1998. He received a PhD in engineering from the same university for his research on digital projection system architectures. In 2004, he joined the Brussels Photonics Team at Vrije Universiteit Brussel, where he was responsible for a research unit that focused on optical modeling and design. He worked on solar concentrators, displays, light-emitting diode (LED)/laser illumination and smart imaging systems, often in close collaboration with SMEs and large-scale companies like BARCO and Philips. In 2013, he started working at the Faculty of Engineering Technology of KU Leuven as an associate professor. His current research interests at the Light and Lighting

Laboratory include the modeling/design of light sources (LED, organic LED, and laser) and advanced illumination systems.

**Geert Deconinck** is a full professor at KU Leuven-Department of Electrical Engineering and has a PhD in applied science. His research interest focuses on control and information technology for power systems. Examples include smart grids, lighting technology, or distributed control of distributed energy resources. He teaches topics related to industrial automation, smart distribution systems, and lighting and measurement techniques. He heads the research group Electrical Energy and Computing Architectures (ESAT/ELECTA) and is a fellow of the Institute of Engineering and Technology and a senior member of the Institute of Electrical and Electronics Engineers.

**Peter Hanselaer** is a full professor at KU Leuven and received his PhD in solid-state physics from the University of Gent (B) in 1986. In 1997, he founded the Light and Lighting Laboratory. His main research areas are lighting, perception and appearance, optical design, light sources and optical metrology. He was recently appointed as editor of Commission Internationale de l'éclairage Division 1, as chair of the Doctoral Committee of the Faculty of Engineering Technology of KU Leuven and as representative of the KU Leuven/ESAT Technology Cluster. He is teaching master courses in photonics, optical fibers, and lighting.