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Abstract. The physical changes of tissue are complicated to evaluate during optical clearing (OC) treatment. Monitoring the changes of optical parameters, including the complex refractive index (CRI), helps people better understand the OC process. From the imaginary part of CRI, we can deduce the extinction coefficient of tissue. Based on the total internal reflection method, the time-dependent CRI of porcine muscle during natural dehydration is well determined. Results show that the real RI increases continuously with the increase of dehydration time, whereas the extinction coefficient initially increases and then decreases. Finally, the extinction coefficient becomes much smaller than the initial value, which demonstrates that better tissue optical clarity is obtained. The change tendency of the extinction coefficient of tissue is used to qualitatively explain the dynamic change of transmittance of a natural dehydrated tissue. Consequently, CRI, especially its imaginary part, is a very useful optical parameter by which to evaluate the OC effect. © 2015 Society of Photo-Optical Instrumentation Engineers (SPIE) [DOI: 10.1117/1.JBO.20.4.045001]

Keywords: optical clearing; tissue; complex refractive index; total internal reflection method; natural dehydration.

1 Introduction

With the development of biological optics, optical diagnostics and optical treatments have drawn much attention. Due to the high scattering of tissue, the penetration depth of light in tissue is limited. Tuchin et al.1 proposed the optical clearing (OC) technique to reduce tissue scattering. The OC technique reduces the mismatches of refractive index (RI) in tissue by partly replacing the interstitial or intracellular water with high RI chemical agents. The RI mismatches are mainly due to the different tissue components. For example, the RIs of collagen fibrils, cellular organelles, and tissue fluid are 1.47, 1.39 to 1.42, and 1.35, respectively.2,3 Rylander et al.5 suggest that water transport is also an important mechanism for tissue OC. According to the heuristic particle-interaction model described in Refs. 4 and 5, the tissue scattering can be controlled by changing the volume fraction of the high RI components in tissue. As studied in Refs. 4 and 6, both the evaporation induced and the compression induced dehydration can allow for better tissue optical clarity.

Various methods have been applied to evaluate the OC effect, such as photographic imaging technique,4 optical coherent tomography technique,4 transmission electron microscopy,4 optical transmittance and reflectance measurements.4,5–10 These methods mostly evaluate the OC effect by detecting the intensity change of light rather than the change of the optical parameter. The inverse adding-double method is a useful approach with which to deduce the absorption and scattering coefficients.11,12 However, the RI value used in the calculation is a supposedly fixed value, which is actually varies during OC. Among the optical parameters of tissue, the complex refractive index (CRI) is an important one and is defined as $n = n_r (1 + i k)$.13 Here, $n_r$ is the real RI, and the imaginary RI, $k$, satisfies $k = \mu_a + \mu_s$. The extinction coefficient, $\mu_e$, represents the energy loss per unit in certain direction caused by absorption and scattering, and $\mu_e = \mu_a + \mu_s$. The total internal reflection (TIR) method is widely used to determine the $n_r$ of tissue15–17 and has been improved to measure the CRI of tissue.18–20 However, the TIR method and CRI have not been applied to evaluate the OC effect.

In this paper, the TIR method is used to monitor the CRI change of natural dehydrated porcine muscle for the first time. A better tissue optical clarity induced by dehydration is observed. However, the change of CRI during dehydration is not monotonic. The heuristic particle-interaction model is utilized to qualitatively explain the change tendency of $k$. Transmittance measurement of the natural dehydrated tissue is also performed, and the dynamic change of the transmittance is well explained by the change tendency of $k$ obtained from the CRI measurement.

2 Materials and Methods

2.1 Materials

Porcine muscle was selected as the sample. The fresh keeping cabinet provides a temperature of −3 to 0°C. The tissue was sealed by preservative film to maintain freshness until it reached room temperature. Before measurement, the tissue was cut into...
50 × 30 mm$^2$ with about a 2-mm thickness. Then the tissue was carefully attached to the prism for CRI measurement and to a slide glass for transmittance measurement. The air gaps between the interfaces should be avoided.

### 2.2 Method

Figure 1 is the schematic diagram of the CRI measurement setup, which is similar to that in Ref. 19. After passing through a beam splitter $M$, a half-wave plate $H$, a polarizer $P$, and an aperture diaphragm $D1$, the $P$-polarized He–Ne laser (632.8 nm) propagates into the prism and irradiates on the tissue.

$D2$ is also an aperture diaphragm. $n_r$ is the real RI of the equilateral triangle prism ($n_p = 1.6166$ at 632.8 nm). $\alpha$, $\beta$, and $\theta$ are the incident angle at the air–prism interface, the apex angle of the prism and the incident angle at the prism–sample interface, respectively. A detector PD1 is used to monitor the power shift of the laser. Detector PD2 is used to detect the emergent light from the prism. The prism is mounted on a rotation stage (PI, M-038), which is controlled by a Mercury C-863 servo motor controller. $\theta$, $\beta$, $\alpha$, and $n_p$ satisfy

\[
\theta = \beta + \arcsin(\sin(\alpha)/n_p).
\]

Based on the Fresnel equation, the reflectance $R_p$ at the prism–sample interface for the $P$-polarized wave is given by

\[
R_p = \left(\frac{n_r^2(1-k^2)\cos\theta - n_p u_2}{n_r^2(1-k^2)\cos\theta + n_p u_2} + \frac{2n_r^2k\cos\theta - n_p v_2}{2n_r^2k\cos\theta + n_p v_2}\right)^2,
\]

where $n_r$ and $k$ are the real part and the imaginary parts of the CRI of sample, respectively. $u_2$ and $v_2$ are two intermediate variables that satisfy

\[
2u_2 = n_r^2(1-k^2) - n_p^2\sin^2\theta
+ \sqrt{[n_r^2(1-k^2) - n_p^2\sin^2\theta]^2 + 4n_p^2k^2},
\]

\[
2v_2 = -[n_r^2(1-k^2) - n_p^2\sin^2\theta]
+ \sqrt{[n_r^2(1-k^2) - n_p^2\sin^2\theta]^2 + 4n_p^2k^2}.
\]

Then the reflectance at the prism–sample interface can be calculated as

\[
R_m = I/I_0(1 - R_i) \quad (5)
\]

\[
R_i = \left(\frac{n_p\cos\alpha - \cos(\arcsin(\sin(\alpha)/n_p))}{n_p\cos\alpha + \cos(\arcsin(\sin(\alpha)/n_p))}\right)^2 (6)
\]

where $I$ and $I_0$ are the light powers of the emergent light and incident light, respectively, and $R_i$ is the power loss at the air–prism interface or prism–air interface.

A nonlinear fitting program based on the Nelder–Mead simplex method is used to simultaneously calculate the $n_r$ and $k$. For the tissue sample, the $n_r$ is the average or effective RI of the tissue components. The inconsistency between the measured data and fitting curve is described by $E^2$, defined as $E^2 = 1 - \Sigma_{m=1}^{n_m}(R_{m;i} - R_i)^2/\Sigma_{m=1}^{n_m}(R_{m;i} - \bar{R})^2$, where $R_{m;i}$ is the $i$th measured reflectance calculated by Eq. (5), $R_i$ is the $i$th calculated reflectance, and $\bar{R}$ is the mean value of measured reflectance over $N$ values of the incident angle. The value of $E^2$ ranges from 0 to 1 and it is closer to 1 when a reliable fitting is obtained.

The reflectance curve of porcine muscle was measured at time points of 1, 2, 3, 4, 5, 6, 7, 8, and 10 h during dehydration. Similar measurements were repeated for three times.

The schematic diagram of the transmittance measurement setup is shown in Fig. 2. A He–Ne laser is used as the incident light. PD1 is used as a monitor to detect the reflected beam from the beam splitter $M$. The transmitted light vertically passes through the slide glass and the tissue and then is collected by an integrating sphere. PD2 is used to detect the emergent light from the integrating sphere. Transmittance measurement was also performed to investigate the OC effect of natural dehydration. The transmittance of the dehydrated porcine muscle was recorded with a time interval of 1 min for a total of 7 h. Similar measurements were repeated for three times.

### 3 Results and Discussion

#### 3.1 Results

The measured reflectance curves of natural dehydrated porcine muscle are shown in Fig. 3. The reflectance curve of the porcine muscle shifts toward the right direction with the increase in dehydration time. Based on the nonlinear fitting program mentioned above, all curves are well fitted with $E^2 > 0.996$ and the CRI values are determined, as demonstrated in Fig. 4. The $n_r$ increases continuously from 1.351 to 1.376 during a 10 h measurement, while $k$ does not increase continuously. During the first stage of natural dehydration (from 0 to 4 h), $k$ increases from 0.0015 to 0.0018. During the second stage
3.2 Discussion

The increase of $n_r$ of porcine muscle during natural dehydration is observed, as shown in Fig. 4. A similar phenomenon can be found in Refs. 22 and 23. The optical parameters of tissue are known to depend on the water content, and both the natural dehydration and compression-induced dehydration can lead to an increase of $n_r$. For the fresh porcine muscle, the measured $n_r$ is 1.351, which is smaller than the value of 1.367 provided in Ref. 19. This might be because the tissue fluid exuded out from the damaged tissue and increased the fraction of fluid at the prism–tissue interface. In our early study, the tissue fluid may lead to a hump in the reflection curve. However, the hump is not observed in this study, as shown in Fig. 3. The main reason may be the difference of tissue preparation. In this study, we used fresh porcine muscle without a freezing treatment, and high pressure was avoided when we attached the tissue to the prism.

It can be seen from Fig. 5 that the growth rate of transmittance is slow during the first few hours and then becomes faster. A similar phenomenon can be found in Ref. 4. From the Bouguer–Beer–Lambert law, collimation transmittance in tissue can be calculated by $A = \exp(-\mu tL)$, where $L$ is the thickness. Thus, both the thickness and the extinction coefficient influence the transmittance. The tissue will shrink due to the water loss and internal structure change, and the shrinkage will increase the transmittance. After 7 h, the thickness of the muscle reduced to about 74%. Considering the change tendency of $k$ observed in this study, we can deduce that the effects of thickness and $k$ on transmittance are opposite during the first stage and then become the same in the second stage. The opposite effects slow down the growth rate of transmittance in the first stage.

In order to explain the change tendency of $k$ during dehydration, the heuristic particle-interaction model is applied. For a dense distribution of scattering particles, the reduced scattering coefficient ($\mu'_s$) is related to the reduced scattering cross section ($\sigma'_s$/cm$^2$) by $\mu'_s = \phi(1-\phi)\sigma'_s/V$, where $\phi$ is the volume fraction of the scattering particles and $V$ is the volume of a single scattering particle. The parabolic relationship between $\mu'_s$ and $\phi$ is demonstrated in Fig. 6. Therefore, the volume fraction of scatters in tissue during dehydration needs to be estimated.
Five porcine muscle samples were weighed and then put into a temperature humidity chamber at 40°C for 72 h to obtain fully dehydrated samples. The weight fractions with an average of about 28% are calculated, as listed in Table 1. According to Ref. 27, the density of dry matter from bovine muscle is $\rho (g/cm^3) = 1.551 - 0.349 \times 10^{-7} T$, where $T$ is the absolute temperature of the sample (measured in Kelvin). We use the data at 20°C (293 K) to estimate the density of fully dehydrated porcine muscle, and the volume fraction of water is deduced as $2.673 \times 1.682 = 27.2$.

<table>
<thead>
<tr>
<th>Sample number</th>
<th>Fresh (g)</th>
<th>Fully dehydrated (g)</th>
<th>Weight fraction (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>8.117</td>
<td>2.250</td>
<td>27.7</td>
</tr>
<tr>
<td>2</td>
<td>6.173</td>
<td>1.682</td>
<td>27.2</td>
</tr>
<tr>
<td>3</td>
<td>16.462</td>
<td>4.666</td>
<td>28.3</td>
</tr>
<tr>
<td>4</td>
<td>10.402</td>
<td>2.959</td>
<td>28.4</td>
</tr>
<tr>
<td>5</td>
<td>9.802</td>
<td>2.800</td>
<td>28.6</td>
</tr>
</tbody>
</table>

The error of $k$ is highly related to the power fluctuation. After being calibrated, the power fluctuation measured by PD2 is less than 0.5%, which corresponds to an error of 0.0001 of $k$.

4 Conclusion

Our experiment uses the TIR method as a useful approach to evaluate the OC effect and shows that CRI is an important optical parameter to understand the OC process. The results also indicate that natural dehydration can provide better tissue optical clarity, while the change of CRI is complicated. The real part of the CRI of porcine muscle increases continuously during the dehydration time, but the imaginary RI initially increases and then decreases. After enough time, the imaginary part of the CRI becomes much smaller than the initial value, revealing that a better tissue optical clarity is obtained. Further studies about OC are expected.

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