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Abstract. We proposed a zoom homogenizer that can control the size of the illumination field by adding one lens array to the conventional imaging-type beam homogenizer. An equivalent lens system was used to derive the imaging condition and size of the illumination field. The result of the ray-tracing simulation shows the validity of the zoom homogenizer to generate the uniform and sharp-edged beam. We presented the strengths and weaknesses of the zoom homogenizer in terms of its etendue with the maximum acceptance angle of the system.© The Authors. Published by SPIE under a Creative Commons Attribution 3.0 Unported License. Distribution or reproduction of this work in whole or in part requires full attribution of the original publication, including its DOI. [DOI: 10.1117/1.OE.57.3.035102]

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1 Introduction

A beam homogenizer, which is one of the beam-shaping methods, has been researched and developed. This method to make the uniform and sharp-edged beam has the advantage of high energy efficiency due to having no energy loss theoretically. The output beam profile of the beam homogenizer is not significantly affected by the input beam condition such as size, shape, and distribution. Recently, the optimization research of the beam homogenizer has been performed to improve the output beam quality.

The zoom function is also useful in many industrial applications. The zoom function changes the irradiated image size so the energy flux or radiant intensity can be easily adjusted. The real-time control of these will be practical or useful in some situations. The conventional beam homogenizer can change the image size by moving the position of the lens array (LA). However, a sharp-edged beam cannot be obtained at all zoom positions.

There have been various studies by adding zoom function in beam homogenization technique to overcome this problem. Researchers have tried to realize the zoom function by modifying the conventional beam homogenizer itself or by adding the lenses to condenser lens part. Several researchers have added a zoom lens to the beam homogenizer and called it a zoom homogenizer. In general, zoom lens systems produce a clear image with a fixed image plane at all zoom positions. But they cannot generate a sharp-edged beam also the position of the image plane varies according to the change in the zoom positions.

In this study, we proposed a zoom homogenizer that can produce a sharp-edged beam at all zoom positions and have a fixed image plane. It was designed by adding an LA to the conventional beam homogenizer. The optical system was assumed as thin lenses, and the imaging condition was derived with equivalent focal length. We confirmed the validity of the proposed method by performing ray-tracing simulation. We also analyzed the maximum angle of the incident beam that the optical system could allow and presented the advantages and limitations of this system regarding the etendue.

2 Derivation of Illumination Field Size According to Zoom Positions

The conventional beam homogenizer can be divided into two types: nonimaging and imaging. The nonimaging type consists of a single LA and a single condenser lens (CL) [Fig. 1(a)]. As there is no imaging condition, a uniform but unclear image is acquired. In contrast, the imaging-type beam homogenizer consists of a pair of LAs and a CL [Fig. 1(b)]. The LA1 plays the role of an object whose image is relayed by the LA2 in the imaging condition. Note that marginal rays are only in Fig. 1(b), because the LA2 plays an aperture stop role in the imaging type. A uniform and sharp-edged image can be obtained at the image plane using this homogenizer type. In this study, we designed an imaging and zooming homogenizer by adding one LA to the conventional imaging-type beam homogenizer.

The zoom optical system implies a system in which the effective focal length (EFL) or the magnification continuously changes while the image plane is maintained. In the proposed optical system, the EFL (f_LA) of LAs is changed continuously, and the image plane is fixed to the focal length (f_c) of the CL. Figure 2(a) shows a zoom homogenizer with three LAs and one CL. For the collimated incident beam, the image plane is at a distance from f_c.

The use of an equivalent lens enables a complex optical system to be simplified using a lens approximation considered as a thin lens. The design using the equivalent lens is useful to derive the size of the illumination field of the imaging beam homogenizer at all zoom positions. Similar to that of the conventional imaging beam homogenizer, the size (D) of the illumination field of the zoom homogenizer is expressed as the ratio of the focal length (f_c) of CL and the EFL (f_LA) of LAs; the pitch size (p) of the LA is as follows:

\[
\frac{D}{p} = \frac{f_c}{f_{LA}}.
\]

\[\text{(1)}\]
In this case, the EFL of three LAs can be written as
\[ f_{\text{LA}} = \frac{f_1 f_2 f_3}{f_1 f_2 + f_1 f_3 + f_2 f_3 - d_{12}(f_2 + f_3 - d_{23}) - d_{23}(f_1 + f_2)}. \]  

(2)

where \( f_1, f_2, \) and \( f_3 \) are the focal lengths of LA1, LA2, and LA3, respectively, and \( d_{12} \) and \( d_{23} \) are the distances between LA1 and LA2 and LA2 and LA3, respectively.

Although there are several variables to determine the size of the illumination field, not every variable is involved in the imaging condition, under which the images of the LA1 lenslets are delivered and perfectly concentrated to the image plane to generate a sharp-edged beam. The key of the imaging condition in the beam homogenizer is to collimate the marginal rays passing through the edge of LA3, which plays the role of aperture stop in this system, to the CL. In a conventional imaging beam homogenizer, only one variable exists regarding the distance: \( d_{12} \). As the imaging condition is fixed, that is, \( d_{12} = f_2 \), zooming is impossible; however, this system comprises two variables regarding distance: \( d_{12} \) and \( d_{23} \) [Eq. (2)]. The determination of the imaging condition for this system can help in designing the zoom homogenizer.

To determine the imaging condition of the zoom homogenizer, we applied the method of the equivalent lens. The green box in Fig. 2(b) displays the simplified lenslet system of LA2 and LA3 in the green box in Fig. 2(a).

As \( f_2 \) and \( f_3 \) are equivalently converted to \( f_{23} \), a principal plane \( H \) is formed between LA2 and LA3. The distance between LA2 and \( H \), that is, \( d_{2H} \) can be written as
\[ d_{2H} = \frac{f_2 d_{23}}{f_2 + f_3 - d_{23}}, \]  

(3)

which determines the position of the equivalent lens of LA2 and LA3. As the equivalently converted focal length \( f_{23} \) is also expressed as \( f_{23} = d_{12} + d_{2H} \), the substitution of Eqs. (3) and (4) in this equation leads to the derivation of the imaging condition
\[ d_{12} = \frac{f_3(f_3 - d_{23})}{f_2 + f_3 - d_{23}}, \]  

(5)

or
\[ d_{23} = \frac{f_2 f_3 - d_{12}(f_2 + f_3)}{f_2 - d_{12}}. \]  

(6)

Substituting Eq. (5) or Eq. (6) in Eqs. (1) and (2), we obtain the size of the illumination field in the imaging condition.
\[ D = \frac{pf_c f_2}{f_3(f_2 - d_{12})} = \frac{pf_c(f_2 + f_3 - d_{23})}{f_2 f_3}. \] (7)

Unlike the conventional beam homogenizer, the distance variables of the lenses are contained and coupled in the imaging condition of the zoom homogenizer. The size of the illumination field is simplified based on the imaging condition. Note that focal length \( f_1 \) of LA1 illustrated as the dotted line in Fig. 2(b) is not involved in the imaging condition [Eq. (5) or Eq. (6)] as well as the size of the illumination field [Eq. (7)]. This is because, as mentioned earlier, LA1 plays the role of an object in the imaging condition.

A zoom ratio is the ratio of maximum EFL to the minimum EFL. \(^15\) The zoom ratio \( R_z \), is defined as

\[ R_z = \frac{f_2 - d_{12,z3}}{f_2 - d_{12,z1}} = \frac{f_2 + f_3 - d_{23,z1}}{f_2 + f_3 - d_{23,z3}} \] (8)

where subscripts \( z1 \) and \( z3 \) are the first and third zoom positions, respectively. To obtain high zoom ratio, it is advantageous to utilize small \( f_2 \) and \( f_3 \) or the highest difference between the distances of the first and third zoom positions.

3 Design Example of the Zoom Homogenizer in Case of \( f_1 = f_2 = f_3 \)

3.1 Ray-Tracing Simulation

In Sec. 2, we derived the size of the zoom homogenizer in the imaging condition using the equivalent lens. This section also shows the result of the ray-tracing simulation conducted using the lens design program CODE V. Because the shape of the lenslets produces the shape of the illumination field at the image plane,\(^7\) we have configured zoom homogenizer with the LAs, which have square-shaped lenslets in the simulation as a design example. Figure 3 shows the lens viewing and spot diagrams in CODE V according to the zoom positions. The simulation parameters are shown in Table 1. For the simple simulation, \( f_1 \), \( f_2 \), and \( f_3 \) are all the same, that is, \( 60.0 \) mm. The distances between LAs, that is, \( d_{12} \) and \( d_{23} \), satisfy the imaging condition. The zoom homogenizer was arranged as an inner zoom, in which the total length of the optical system is fixed (Fig. 3). LA2 and LA3 play the role of variator and compensator, respectively. The blue line connected to LAs according to the zoom position present the zoom locus. The position of the image plane is fixed even though \( f_{LA} \) changes. However, the image size \( (D) \) of the illumination field changes, thus satisfying the image condition, Eq. (7), according to the zoom positions. Figure 3 shows a spot diagram at each zoom position with a sharp-edged square-shaped beam. According to the beam size (15.0 mm) at the second zoom position, when \( f_{LA} \) is minimum, the beam size is enlarged to 18.3 mm, and when \( f_{LA} \) is maximum, it reduced to 11.7 mm. The simulation result shows that the imaging condition in the zoom homogenizer is reliable. Furthermore, the zoom ratio is calculated using Eq. (8) and is ~1.6.

Figure 4 shows the magnification and normalized energy fluence according to the change of \( d_{12} \). The magnification is defined as the ratio of the size \( (D) \) of the illumination field at

\[ \begin{array}{cccccc}
\text{Zoom1} & 27.3 & 10.0 & 32.7 & 32.7 & 32.7 & 18.3 \\
\text{Zoom2} & 20.0 & 30.0 & 20.0 & 200.0 & 40.0 & 15.0 \\
\text{Zoom3} & 8.6 & 50.0 & 11.4 & 11.7 & 11.4 & 11.7 \\
\end{array} \]

Table 1 Parameters of the example simulation in (mm).

![Fig. 3](https://www.spiedigitallibrary.org/journals/Optical-Engineering/035102-3/March-2018-Vol.57(3)/Kim-et-al.-Design-of-zoom-homogenizer-to-control-size.png)

Fig. 3 Lens viewing and spot diagram of the design example of the imaging condition according to the zoom positions. The blue solid lines represent the zoom locus of the variator (LA2) and compensator (LA3).
each zoom position to the size \((D_{z2})\) of the second zoom position. The maximum and minimum magnifications are \(\sim 1.2\) and 0.8 at zoom 1 and zoom 3, respectively. The energy fluence, beam intensity, or radiated area are a more direct and important factor than the image size in the industrial field. As aforementioned, the normalized fluence is the ratio of energy fluence \((J)\) at each zoom position to the energy fluence \((J_{z2})\) at the second zoom position. The maximum and minimum values of the normalized fluence are \(\sim 1.7\) and 0.7, respectively.

### 3.2 Pros and Cons of the Proposed System

In Secs. 2 and 3.1, we have assumed that the incident beam is parallel to the zoom homogenizer. But the incident beam may not be collimated or parallel to the zoom homogenizer in practical situations. The beamlets passing through the lenslet create a crosstalk that invades the adjacent area of the next LA. It happens when the incident beam has an angle that exceeds the etendue of the system.\(^{30,31}\) The etendue of the light source should be lower than the etendue of the system to prevent this kind of situation. We have analyzed the etendue of the system in the maximum acceptance angle, which occurs with no overflow at all zoom positions. The following system matrix gives the information of the angle and the image size at the image plane when the incident ray has an arbitrary angle.

\[
\begin{pmatrix}
\theta_{img} \\
h_{img}
\end{pmatrix} = \begin{pmatrix}
1 & 0 \\
f_c & 1
\end{pmatrix} \begin{pmatrix}
1 - \frac{f_c}{f} & 1 \\
0 & 1
\end{pmatrix} \begin{pmatrix}
1 - \frac{f_c}{f} & 0 \\
0 & 1
\end{pmatrix} \begin{pmatrix}
d_{12} & 0 \\
0 & 1
\end{pmatrix} \begin{pmatrix}
\theta_{LA1} \\
h_{LA1}
\end{pmatrix}.
\]

In Eq. (9), \(\theta_{LA1}\) and \(h_{LA1}\) are the angle and the height of the incident ray at LA1, respectively, and \(\theta_{img}\) and \(h_{img}\) are the angle and the height of the image at the image plane, respectively. A local coordinator of the \(n\)'th array \(\Delta h = -np\) determines the position of the rays according to the lenslet channel. The ray height at the image plane, \(h_{img}\), in imaging condition is

\[
h_{img} = \frac{pf_c f_2}{2f_3(d_{12} - f_2)}.
\]

and is a half of the image size \((D)\) of Eq. (7). It is worth noting that \(h_{img}\) is not influenced by the incident angle to LA1, \(\theta_{LA1}\). It is valid in the thin lens approximation scheme.

We have to consider the cases of ray clipping by LA2 or LA3 to determine the maximum acceptance angle for the zoom homogenizer. In this system, two cases are significant (Fig. 5). The first case \(\theta_{max 1}\) is that the incident ray, which has an angle starting from the upper edge \((p/2)\) of LA1 hits the upper edge of the LA2. In this case, the angle corresponds to the numerical aperture of LAs \((\text{NA}_{LA})\) in common with the conventional beam homogenizer.\(^{31}\) Thus, it is a fixed value regardless of the zoom positions (red solid line). The second case \(\theta_{max 2}\) is that the incident ray hits the lower edge \((-p/2)\) of the LA3. The maximum angle varies as the zoom position changes in this case (blue solid line). These two cases are dominant, and they correspond to the maximum acceptance angles in this system. The other cases are not meaningful in this system.

The dashed lines in Fig. 5 mean mirrored values of each solid line, respectively, by considering the symmetry of the optic axis of the incoming ray. The pairs of the original line and the flipped line form the acceptable angular range. The case that has the smallest absolute angle has to be selected to define the angular limitation for this system among these cases. In other words, it is required to limit the zoom range when the angle of the input source is partially out of acceptable angular range.

The product of the angle \(\theta_{img}\) and the height \(h_{img}\) of the ray is almost constant when \(\theta_{LA1} = \theta_{max 1}\), \(n = 5\), and \(d_{12}\) is in zone II. Especially, it can become the etendue-conserving system of the output product \((\theta_{img} h_{img})\), which is the same as the input product \((R \cdot \text{NA}_{LA})\) when \(d_T = f_c\).\(^{21}\)

### 4 Conclusion

We have designed the zoom homogenizer using three LAs, and it is confirmed by using the ray-tracing simulation. The image size is changed continuously with having no variations in its image distance and its edge steepness. The zoom ratio was 1.6 in the design example, which is configured as an inner zoom with the same three LAs. The zoom
homogenizer works well in the acceptable angular range without a crosstalk in the LA system. It will be useful in many industrial applications, which are required to regulate the intensity or energy fluence with a sharp edge steepness.

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References

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