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# Graphene photonics for optical communications

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## ABSTRACT

Graphene is a suitable material for optoelectronic applications that reveals several advantages and complementarities compared with other technologies. Graphene is a gapless material that absorbs radiation from visible to far-infrared and beyond including the terahertz range. The absorption can be modified by changing the material doping, i.e. Fermi level energy, by applying an external electric field. The change in absorption can be very fast and, for this reason, graphene can be used to realize optical modulators. Moreover, the absorption change occurs along with refractive index change. In particular, conditions, i.e. for Fermi level above the Pauli blocking, light is not absorbed anymore and only phase change occurs. Phase change has been demonstrated to be good for fast Mach Zehnder interferometer based modulation. The broadband absorption of Graphene is also exploited to realize efficient photodetectors. Generation of hot electrons upon light absorption in graphene is the cause of photo-thermal effect (PTE) that leads to photo voltage generation. PTE is an ultrafast process that is used for fast photodetection. In this work, the vision for graphene-based integrated photonics is presented. We review state-of-the-art graphene-based modulators and detectors and outline a roadmap matching the technology readiness requirements with the datacom and telecom market demands. We show that graphene-based integrated photonics could enable ultra-high spatial density, low power consumption for board and intra- data centre connectivity, access networks, metropolitan, core, regional and long-haul optical communications.

**Keywords:** Graphene, integrated photonics, optical communications

## 1. INTRODUCTION

Photonics is a key enabling technology that is continuously growing in terms of markets and applications. Notwithstanding the number of applications, photonics production is not yet comparable to microelectronics in terms of volumes even though most of our communications and many other applications rely on photonics [1]. In optical communications, Telecom and even more Datacom, photonics plays a major role; in fact, all the large bandwidth and connection capacity needs to be optical. The main utilization of photonics is in the transport, metro, access network and in data center interconnections. These applications require volumes not exceeding million devices per year at most. However, in the next years this volume size will change because of a transformation of the connection requirement due to the emerging market of IoT (Internet of Things) [2]. The very large increase of connected devices will require a large and pervasive photonic communication infrastructure and ultra large bandwidth in all layer of the networks [3]. In the 5G era, that is under development, the supposed increase of bandwidth will be a factor of 1000 with respect to present 4G bandwidth [4]. This will require large bandwidth photonic interconnections starting from wireless antennas down to base stations, to radio access network and so on. In data centres, the interconnections are in between racks whereas in the near future, the optical interconnections will be extended to intra- board interconnection that will increase the volume of photonic devices [5]. In other words, the number of photonic interconnection will grow to tens of million devices per year.

The photonic devices, e.g. modulators and detectors, need to meet strict requirements in terms of size, efficiency and power consumption. Graphene photonics is a new and promising technology, still under development [6]. Graphene photonics provides all the basic components that are required for any communication system, i.e. switches, modulators and photodetectors. Graphene is a gapless material that exhibits both electro- absorption and electro- refraction which can be exploited for light modulation and photodetection [7]. Graphene photonics is compatible with complementary metal oxide semiconductor (CMOS) processing, and enable post-processing fabrication of active devices, e.g. modulators and detectors. Graphene photonics provides significant advantages with respect to standard Silicon Photonics. Practically, the use of Graphene on optical waveguides means a post- processing shift in manufacturing from front- end to back- end-of-line. Moreover, graphene technology does not necessarily require expensive SOI wafers, or implantation for junctions, and Ge growth for detectors. These features imply a simpler technology at reduced costs.

Here we report on recent progresses of Graphene photonics, providing a review of the latest results about modulators and detectors. Section 2 is dedicated to Graphene modulators, while section 3 is dedicated to Graphene photodetectors.

## 2. GRAPHENE PHOTONICS MODULATORS

Electro-optic modulator is one of the key devices of photonic integrated circuits. Broadband operation at high modulation speed combined with small footprint, low power consumption and minimum insertion loss represent the main challenge for future optical interconnect [8]. Graphene based optical modulators offer several advantages with respect to silicon based modulators: broadband electro-absorption, efficient electro-refractive effect, high speed and athermal optoelectronic properties. Graphene, a zero bandgap material, is a broadband absorber with 2.3% absorption at normal incidence [9]. However, Graphene can be integrated on top of an optical waveguide to enhance the interaction of the propagating guided mode and improving significantly the light absorption. Moreover, the complex dielectric constant of Graphene, i.e. refractive index and absorption, can be properly changed by varying the electrical carrier concentration through electrical gating, allowing for the realization of broadband electro-absorption modulator (EAM) or phase modulators [10]. The first EAM based on a single layer of graphene on top of a doped Si waveguide was demonstrated in 2011 [11]. The Graphene monolayer was grown by chemical vapour deposition (CVD) and transferred on a 7nm thin film of Alumina ( $\text{Al}_2\text{O}_3$ ) covering a doped Si ridge waveguide. The resulting Graphene- $\text{Al}_2\text{O}_3$ -Si capacitor is the active device, where an external voltage applied to the capacitor allows for the gating of the graphene Fermi level energy. This work experimentally showed for the first time the great advantage of Graphene EAM's with respect to the Franz Keldish (FK) and quantum confined stark effect (QCSE) modulators. Typical Si photonics FK and QCSE EAM based on SiGe or III-V on Si are limited by the bandgap of the material requiring different devices for different operating wavelength with a typical bandwidth of 30nm. On the contrary, Graphene allowed broadband operation over hundreds of nanometers from the telecom O band beyond the C band towards the mid and far infrared.

The industry's first Graphene on SOI EAM capable of 10Gb/s non-return-to-zero NRZ modulation was demonstrated in 2016 [12]. In this device, the authors used a single layer of CVD grown graphene separated from the doped Si ridge waveguide by 5nm thin  $\text{SiO}_2$  layer in order to create Graphene- $\text{SiO}_2$ -Si capacitor stack. The authors demonstrated open eye diagrams at 10Gb/s with extinction ratio as high as 2.5dB, a peak-to-peak driving voltage of 2.5V, small foot-print ( $50 \times 10 \mu\text{m}^2$ ) and energy consumption as low as 350fJ/bit. However, this kind of modulators has a strong limitation due to the use of implantation of dopants in the Si waveguide in order to achieve high speed of operations. An improved design is obtained by implementing the gating capacitor by means of two layer of graphene separated by a thin dielectric film. In this way, the optical waveguide can be a simple passive waveguide of any material. The double layer graphene modulator was demonstrated for the first time in 2012 [13]. Recently, two layers of graphene separated by a 65nm thin film of  $\text{Al}_2\text{O}_3$  were used on a  $\text{Si}_3\text{N}_4$  microring resonator to realize a 22Gb/s electro-optic modulator based on resonator loss modulation at critical coupling [14]. The device is based on the Zeno effect induced by tuning the optical absorption of Graphene on top of a Si microring thus changing the critical coupling condition of the resonance. The authors demonstrated 15dB static ER with a record modulation bandwidth of 30GHz.

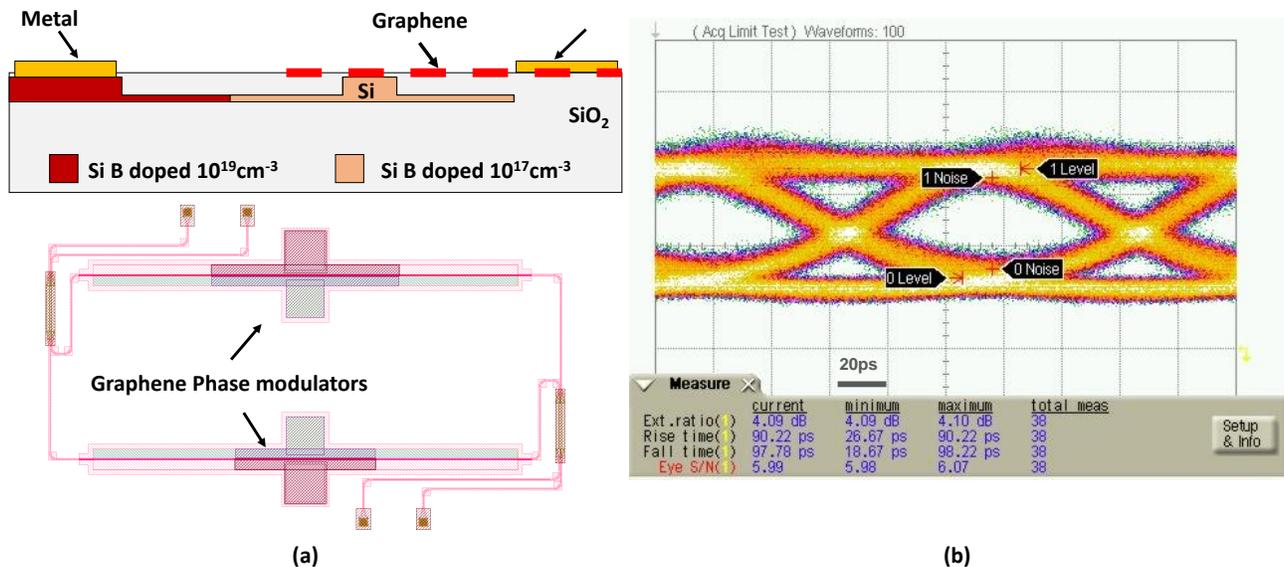


Figure 1. (a) Cross-section of the device and MZI layout. (b) NRZ eye diagram at 10Gb/s

Graphene can be used also for the realization of efficient phase modulators. These are of great interest in order to exploit complex modulation formats as for example QPSK and QAM, or in order to realize MZI modulators with low insertion loss and large extinction ratios. We demonstrated that phase modulation in single layer graphene on Si EAM introduces a chirp that is beneficial in compensating the chromatic dispersion of single mode optical fibres. This effect was exploited to demonstrate optical transmission of a 10Gb/s NRZ data stream up to 100km single mode optical fibre link [15]. Phase modulation has been also demonstrated in in single layer graphene on Si modulators integrated in MZI up to 10Gb/s [16]. Figure 1 shows the cross section of the modulator and the demonstrated results. We also investigated the double layer structure on both Si and Si<sub>3</sub>N<sub>4</sub> waveguides to demonstrate efficient and fast modulators. Figure 2 shows the cross section of the proposed device.

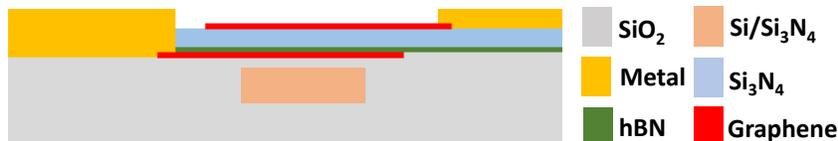


Figure 2. Cross-section of the double layer graphene modulator integrated on Si or Si<sub>3</sub>N<sub>4</sub> waveguides

Graphene has been grown with a predetermined seeded array approach [17] and the resulting single crystal graphene had a defined size that was changed from 100 to 200µm depending on the application. The crystals were then transferred with a semi-dry approach to the waveguides previously planarized with a residual cladding thickness of 10nm on to the top of the core. The flatness of the planarized surface was determined with AFM measurement to be 0.6nm. The single crystal transferred on this surface was analysed with a Raman probe and the carrier mobility resulted to be 7000cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup>. A commercial hBN single layer was transferred on the graphene crystal. The role of hBN is only for protection during the step of 10nm thin layer PECVD Si<sub>3</sub>N<sub>4</sub> deposition that was used as spacer in the capacitor. The second single crystal of graphene was transferred on top of Si<sub>3</sub>N<sub>4</sub>; the full stack was then patterned for metallization. Two Ni/Au contacts were then evaporated on the pad areas. For the SOI waveguide the same stack and contact scheme were used. A number of graphene EA and MZI modulators have been realized on Si<sub>3</sub>N<sub>4</sub> waveguides. The cross section of the Si<sub>3</sub>N<sub>4</sub> waveguide was 250nmx1500nm. Figure 3 shows optical microscope pictures of the fabricated devices and the obtained NRZ eye diagrams. The modulators were electrically driven at 3.5V<sub>pp</sub> with a 2<sup>31</sup>-1 pseudo random binary sequence (PRBS) from the pattern generator at 10Gb/s, 20Gb/s and 25Gb/s.

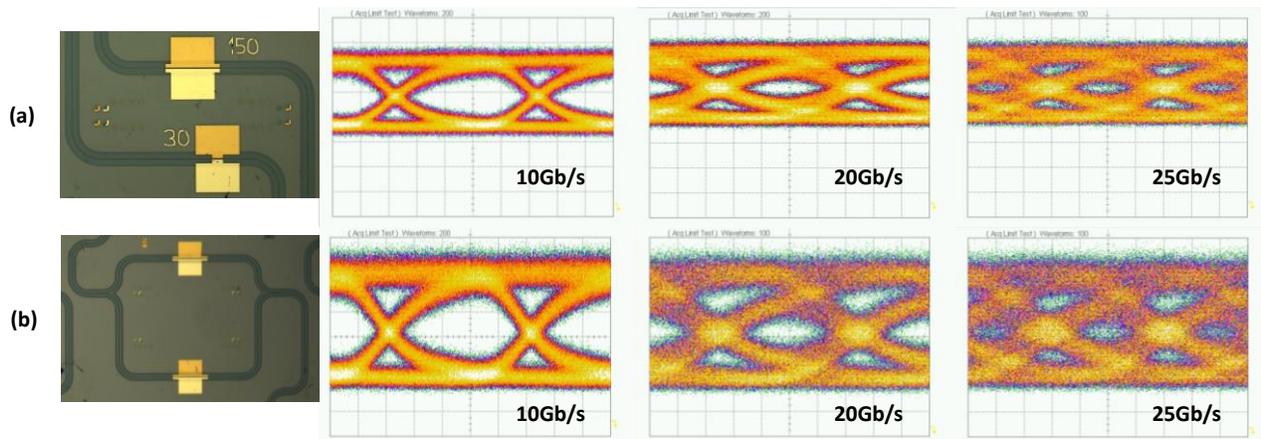


Figure 3. (a) Optical microscope image of a fabricated double layer graphene EAM on Si<sub>3</sub>N<sub>4</sub> waveguide and measured NRZ eye diagrams at 10Gb/s, 20Gb/s and 25Gb/s. (b) Optical microscope image of a fabricated double layer graphene MZI on Si<sub>3</sub>N<sub>4</sub> waveguide and measured NRZ eye diagrams at 10Gb/s, 20Gb/s and 25Gb/s.

We also demonstrated double layer EAM modulator integrated on 220nmx480nm SOI nanowires. In this case, the contacts on graphene are closer because of the smaller cross section of the SOI nanowire compared to the Si<sub>3</sub>N<sub>4</sub> waveguide. The effective -3dB electro-optical bandwidth of the fabricated double layer EAM on SOI is ~30GHz. The modulator was

electrically driven at  $3.5V_{pp}$  with a  $2^{31}-1$  pseudo random binary sequence from the pattern generator at 10, 25 and 50Gb/s. Results are shown in Fig. 4

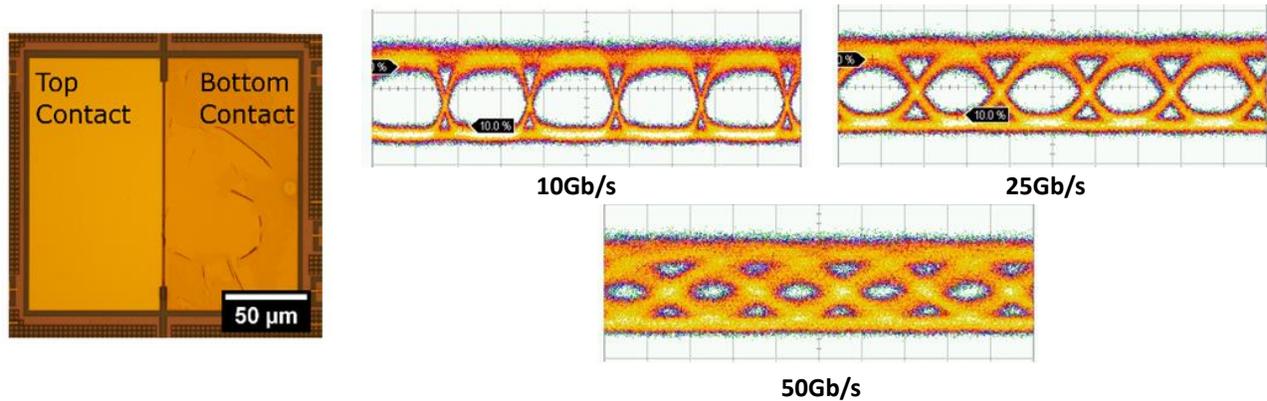


Figure 4. Optical microscope image of a fabricated double layer graphene EAM on SOI waveguide and measured NRZ eye diagrams at 10Gb/s, 25Gb/s and 50Gb/s.

### 3. GRAPHENE PHOTONICS DETECTORS

Graphene photonics photodetectors should be integrated in a waveguide and optimized in performance as photodetectors relevant for Si photonics. Photodetectors in Si photonics are miniaturized and simple to fabricate photodetectors fulfilling the demands of scalability, footprint, and cost. The graphene photodetectors responsivity  $R_{ph}$  should be comparable to that of Ge devices in Si photonics, that is  $\sim 0.85 - 1.15$  A/W at 1,550 nm [18].

Graphene provides significant advantages for photodetection due to its unique properties [19]: wavelength independent absorption, fast photovoltage generation, fast photo-switching rates, high internal quantum efficiency  $>80\%$ . Graphene photodetectors may be based on different detection mechanisms including: photovoltaic (PV) effect [20], photo-thermoelectric (PTE) effect [21] or bolometric effect [22]. Among these, the PTE effect is very promising for the realization of photodetectors with both high-speed (up to 100 GHz) and efficiency that may exceed 100% because of hot carrier multiplication [23]. The PTE effect arises in the presence of photogenerated hot carriers generated on a few tens of femtosecond timescale. The photogenerated hot electrons and holes are thermally decoupled from the crystal lattice because of slower electron-lattice relaxation time, in the picosecond time scale [21]. The temperature profile of the hot carriers generated through the optical absorption may be converted in a local photovoltage via the Seebeck effect occurring when the Seebeck coefficient is changed along the graphene layer. This is the case of double-gated devices where the graphene Seebeck coefficient is changed by applying different voltages to different regions of the graphene layer. The photovoltage generated through the PTE effect exhibits significant advantages, e.g. it can be amplified by a voltage amplifier rather than a transimpedance amplifier that allows the realization of photoreceivers at lower cost and with lower power consumption.

The integration of graphene on Si photonic for the realization of photodetectors have been reported. These devices are typically based on a single layer of graphene evanescently coupled to Si waveguides provided with two metal contacts. The guided mode enables longer interaction between graphene and the optical waveguide compared with normal-incidence illumination, allowing for larger optical absorption and responsivity. A graphene photodetector based on the bolometric effect consisting of a CVD grown single layer of graphene on a Si waveguide operating at 1550 nm has been demonstrated with a -3 dB bandwidth of 41GHz operating at 50Gb/s [24]. A bolometric waveguide-integrated graphene photodetector with plasmonic enhancement has been recently demonstrated with 0.5A/W responsivity at telecom wavelengths and photoresponse faster than 110GHz [25].

A PTE graphene photodetector integrated on a Si waveguide was also demonstrated with a responsivity of 0.36 A/W and 40 GHz bandwidth [26]. The device was obtained by asymmetrically placing the drain and source metal contact with respect to the waveguide core and using a gate to maximize the Seebeck coefficient of the graphene. The single layer graphene was encapsulated in hexagonal boron nitride (hBN) to improve the mobility up to  $\sim 40,000\text{cm}^2\text{V}^{-1}\text{s}^{-1}$ , therefore the Seebeck coefficient. Double gated PTE photodetectors have been also demonstrated. A Si slot-waveguide has been used

as a dual gate for the realization of a graphene PTE photodetector. The device exhibits extrinsic responsivity of 3.5V/W at zero bias and a -3dB bandwidth of 65GHz [27]. A similar approach was recently used with a photonic crystal waveguide where the photonic structure is used for mode confinement and as a split-gate electrode in proximity of the optical absorption region. A photoresponsivity of 4.7 V/W and a (setup-limited) electrical bandwidth of 18 GHz are achieved [28].

We used the split-gate geometry to realize a PTE graphene photodetector on a Si<sub>3</sub>N<sub>4</sub> waveguide. The cross-section of the device is shown in fig. 5.

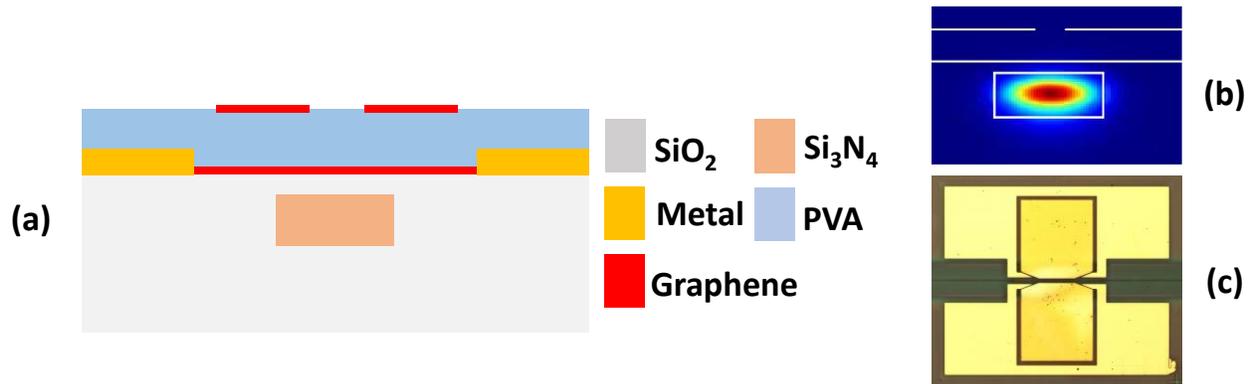


Fig. 5: (a) Cross section of the device, (b) Optical mode of the waveguide, (c) Optical micrograph of a realized device

The device based on CVD Graphene single crystal transferred by the semi-dry method on top of a single mode transverse electric (TE) Si<sub>3</sub>N<sub>4</sub> waveguide with cross-section 1.5μm x 250nm. The photodetector is based on a split gate structure in order to create the two regions with different Seebeck coefficients by electrical gating of the bottom graphene with different chemical potentials. The top gates are obtained with a second layer of CVD graphene transferred with the same semi-dry method, and separated from the bottom layer by 120nm thick of polyvinyl alcohol (PVA). The bottom and top graphene are electrically accessible through Nickel/Gold contacts.

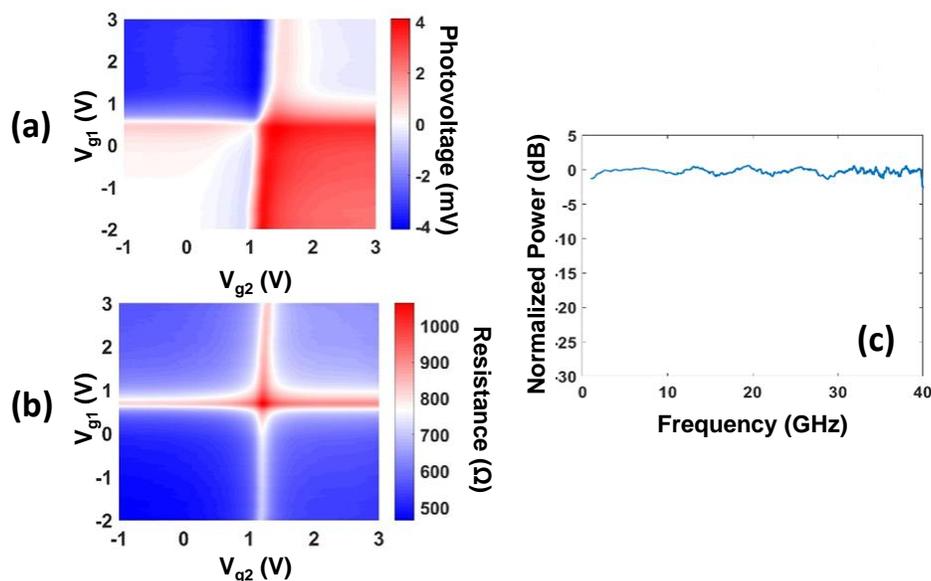


Fig. 6: (a) Photovoltage map as function gate voltages, (b) Resistance map, (c) Measured optical bandwidth of the device

We optimized the cross section of the device by numerical simulations including optical mode analysis at 1550nm (fig. 1(b)). We modelled the PTE effect to describe the case of a detector integrated on a photonic waveguide [29]. The model is defined by the system of equations:

$$\nabla^2 T_{el} + \frac{1}{L_c^2} (T_{el} - T_0) = \frac{P_{abs}(x, z)}{k_{el} h_g}; \quad \vec{j} = -\sigma \nabla V + \sigma S(\mu(x)) \nabla T_{el}; \quad \vec{\nabla} \cdot \vec{j} = 0$$

We characterized the detector at 1550nm wavelength in both static and dynamic operations. We measured the DC photovoltage and photoresistance map sweeping the two gate voltages, as well as the electrical bandwidth of the device by using a Mach Zehnder modulator up to 50GHz and a vector network analyzer (VNA). Figure 6 shows the results.

The measured photovoltage map (fig. 6(a)) exhibits a clear six-folds pattern, which is a clear signature of the PTE effect [29]. The maximum corresponds to a responsivity of 6V/W. The resistance map of the device, shown in fig. 6(b), exhibits a maximum, i.e. Dirac point, close to 0V that indicates that the transferred graphene is a high quality material. This aspect is very important in order to demonstrate a clear PTE effect.

The electrical bandwidth of the device was obtained using a Lithium Niobate Mach Zehnder modulator with 50GHz electro-optical bandwidth and a vector network analyzer (VNA). The gate voltages were set to obtain the maximum responsivity. The setup was calibrated using a commercial photodiode with 40GHz bandwidth. The result shown in fig. 6(c) demonstrates a flat bandwidth up to 40GHz limited only by the experimental setup. This result is very encouraging because obtained with a wafer-scale process with a CVD material.

#### 4. CONCLUSIONS

The results shown in this paper are obtained on a wafer-scale integration process on different photonic platforms employing CVD graphene. We demonstrate that double layer graphene modulators are ready to compare with state-of-the-art photonics integrated devices and paves the way to the establishment of Graphene photonics technology which could enable ultra-high spatial bandwidth density, low power consumption for connectivity within boards or between data centres, access networks and metropolitan, core, regional and long-haul optical communications.

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